

# *Jet finding at the LHC era*

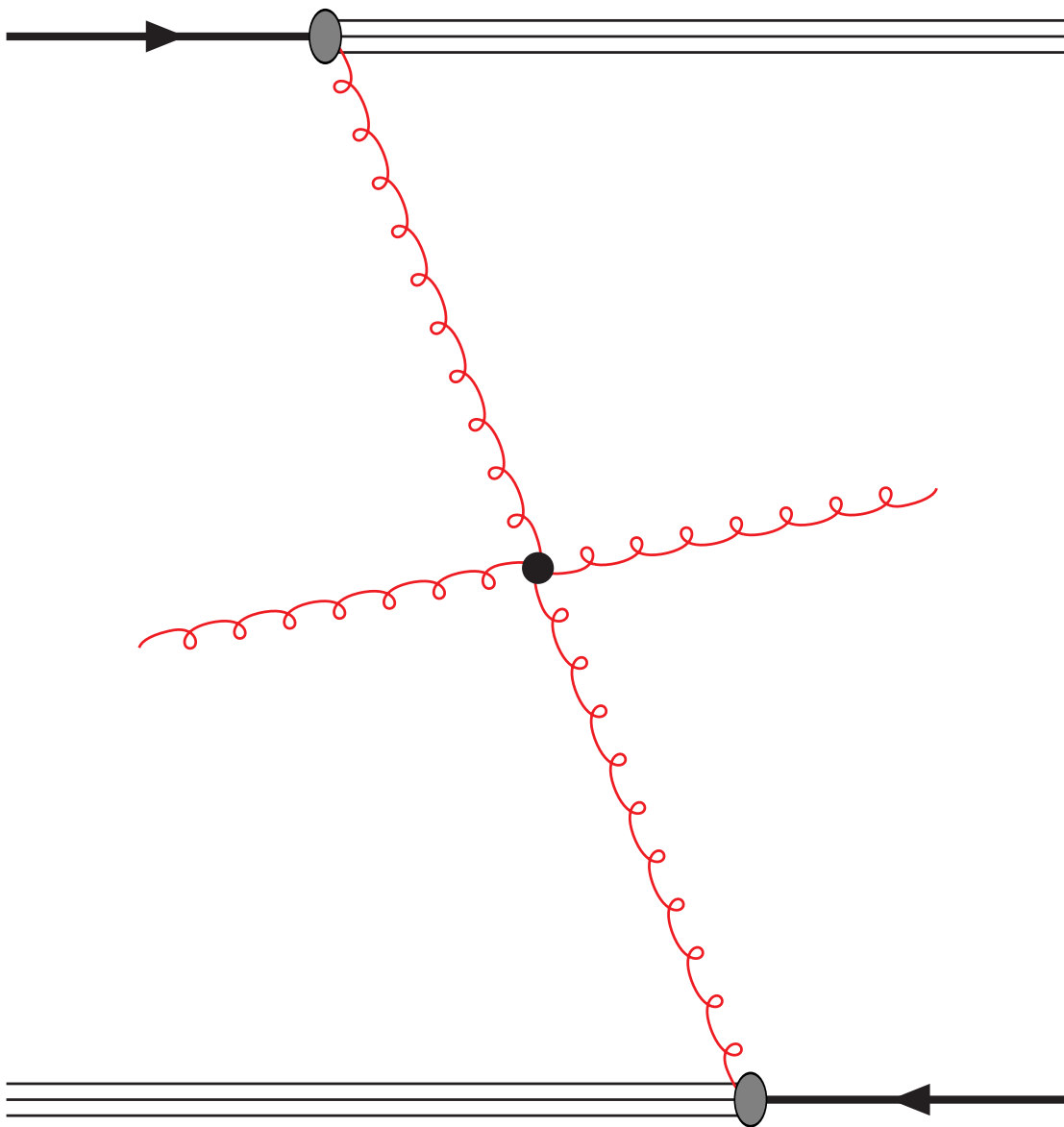
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Brookhaven National Laboratory

in collaboration with G. Salam, M. Cacciari and J. Rojo  
arXiv:0704:0292, arXiv:0802:1188, arXiv:0802:1189,  
arXiv:0803.0678 + works in preparation

- Foreword: why jets? what are they?  
introducing the basic concepts
- Part 0: recent progress: building a solid toolkit  
jet definitions meeting the fundamental requirements
- Part 1: jets in  $pp$  collisions
  - **Choosing the adapted jet definition**  
which jet algorithm is best suited?
  - **Subtracting pileup background using jet areas**
    - defining areas
    - analytic control
    - using them for pileup subtraction
- Part 2: jets in heavy-ion collisions
  - subtraction subtleties
  - preliminary results

# ***Foreword: why jets? what are they?***



## Hard scattering ( $2 \rightarrow n$ )

computed exactly at  $\mathcal{O}(\alpha_s^p)$

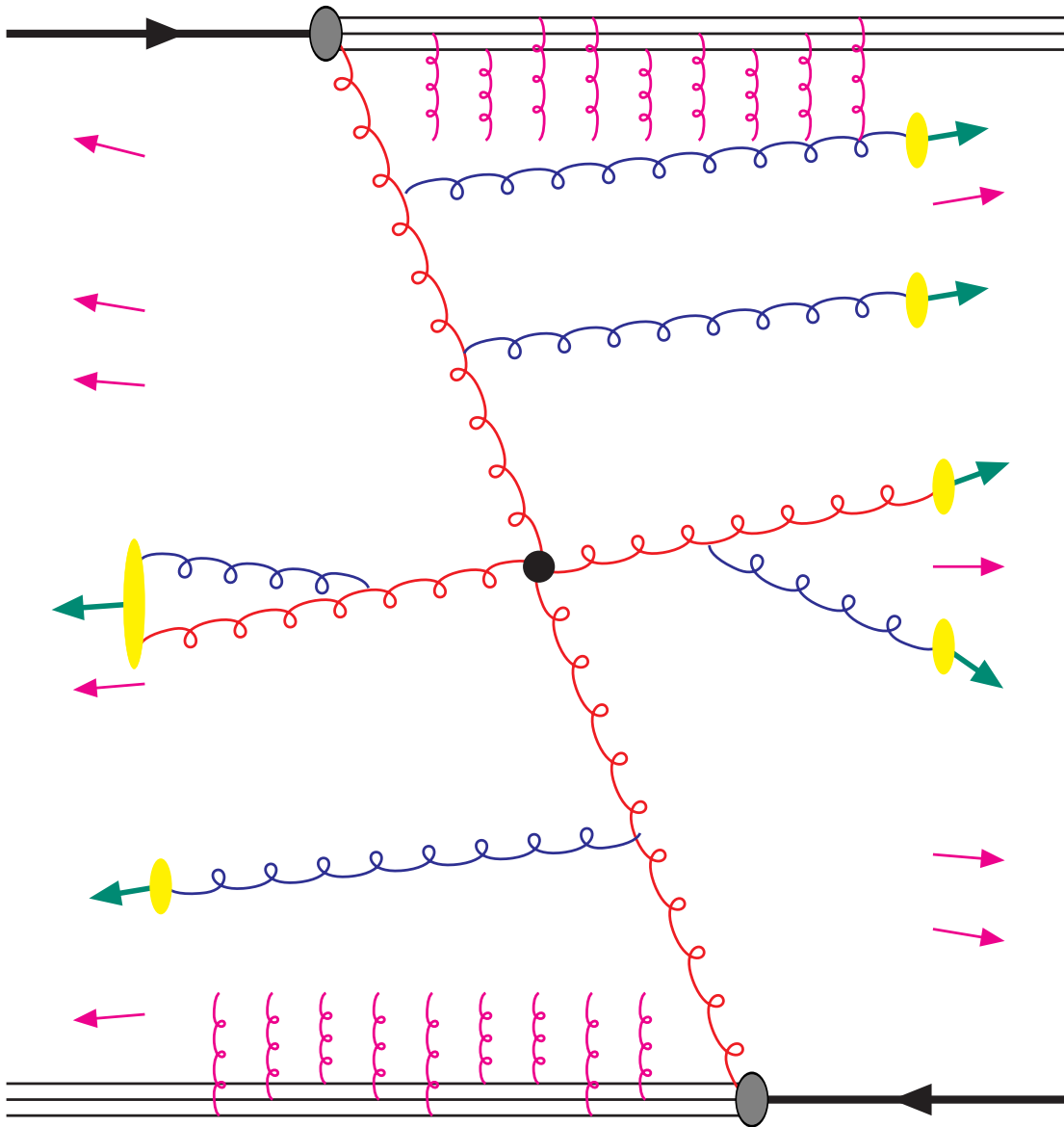
$$gg \rightarrow gg, gg \rightarrow ggg,$$

$$gg \rightarrow gggg,$$

$$gg \rightarrow H \rightarrow b\bar{b},$$

$$gg \rightarrow t\bar{t} \rightarrow \mu\nu_\mu b\bar{b}q\bar{q},$$

$$gg \rightarrow Z' \rightarrow q\bar{q}, \dots$$



Hard scattering ( $2 \rightarrow n$ )

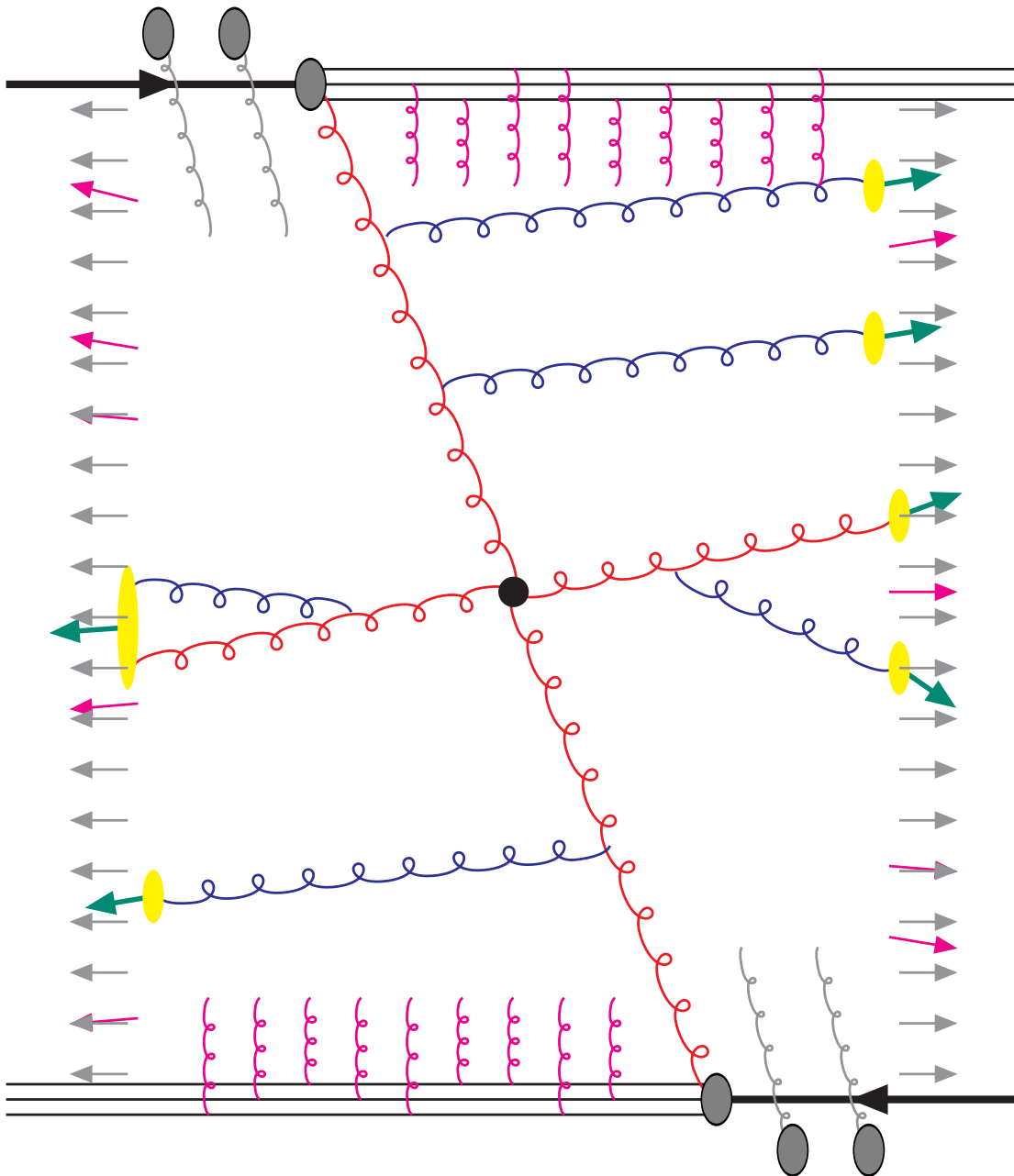
Parton level

$\approx$  resummed collinear div.  
 $\sum_i \alpha_s^i \log^i(p_t^2/\mu^2)$

Hadron level: hadronisation

Underlying event

beam remnants interactions  
 $\Rightarrow$  soft background



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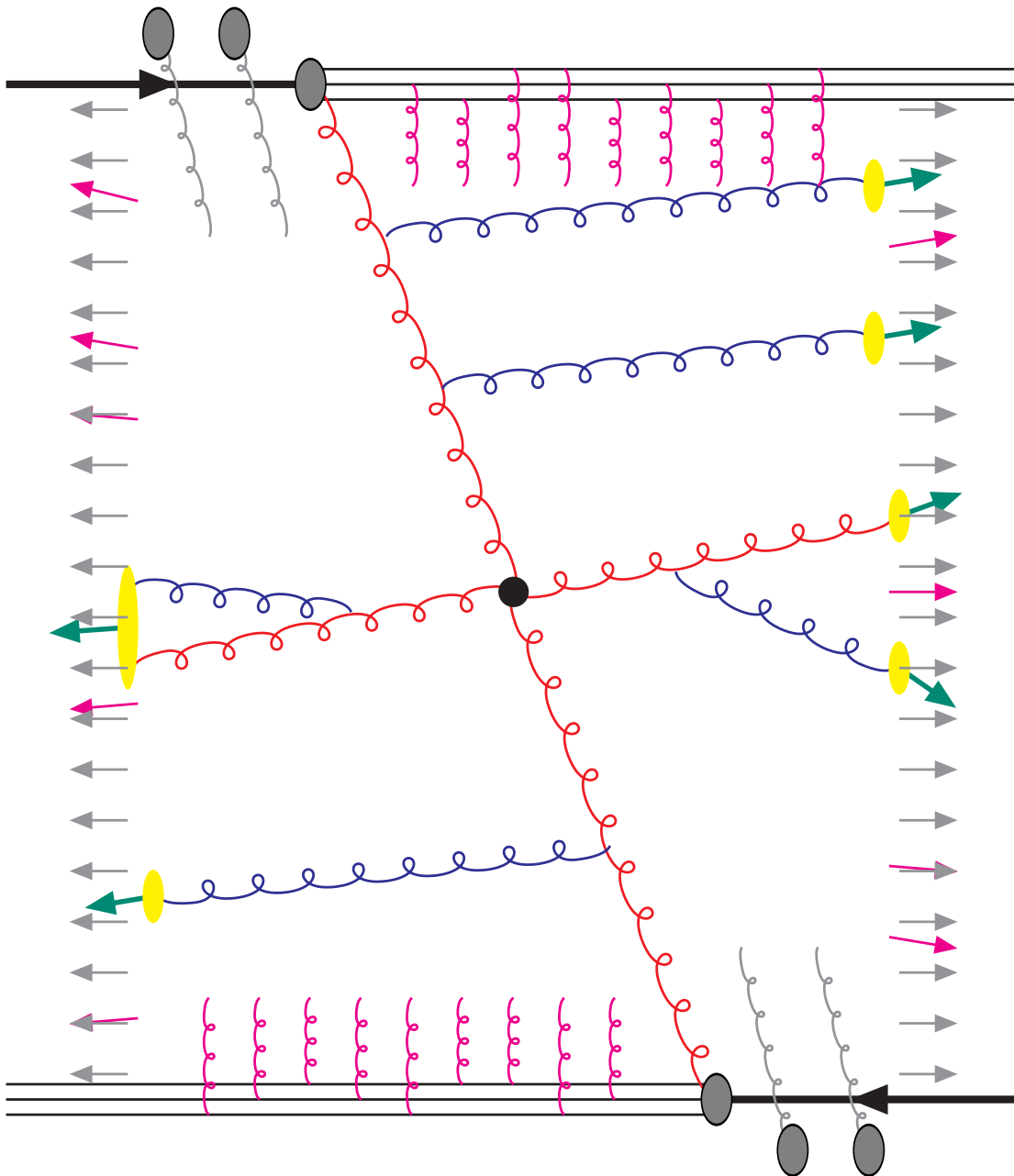
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$\Rightarrow$  soft background

Pileup

$\approx$  uniform soft background



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“Jets”  $\equiv$  hard partons

Parton ambiguous

$\Rightarrow$  multiple jet definitions

Class 1: recombination	Class 2: cone
Successive recombinations of the “closest” <sup>(a)</sup> pair of particle	find directions of energy flow ≡ stable cones <sup>(b)</sup>
Nice perturbative behaviour	Small sensitivity to soft radiation (UE,PU)
Often used in $e^\pm e^\pm, e^\pm p$	Often used in $pp$

(a) Distance: (stop when  $d_{\min} > R$ )

$$k_t: \quad d_{i,j} = \min(k_{t,i}^2, k_{t,j}^2)(\Delta\phi_{i,j}^2 + \Delta y_{i,j}^2)$$

$$\text{Aachen/Cam.}: \quad d_{i,j} = \Delta\phi_{i,j}^2 + \Delta y_{i,j}^2$$

(b) stable cones (radius  $R$ ) such that:

the total momentum of its contents points in the direction of its centre



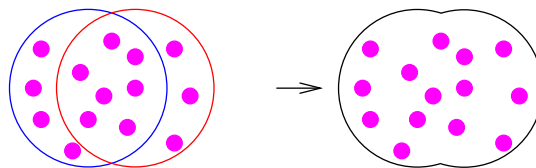
- Seeded (iterative) approaches: iterate from an initial position until stable
    - seed = initial particle
    - seed = midpoint between stable cones found at first step
  - One has to deal with overlapping stable cones: 2 subclasses
-

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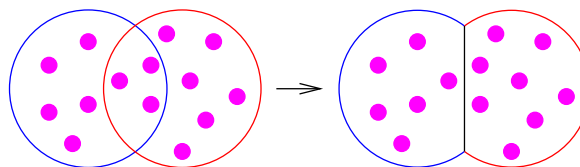
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Class 2(a): cone with split-merge (ex.: JetClu, Atlas, MidPoint):

$$\tilde{p}_{t,\text{shared}} > f\tilde{p}_{t,\text{min}}$$



$$\tilde{p}_{t,\text{shared}} \leq f\tilde{p}_{t,\text{min}}$$

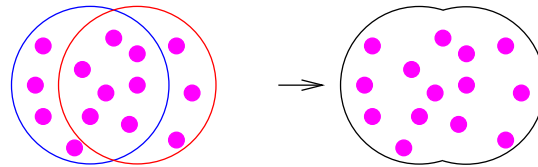


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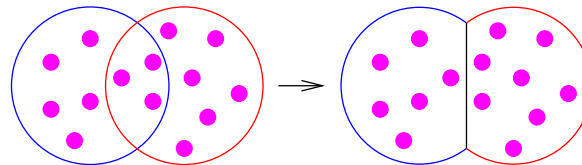
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Class 2(b): cone with progressive removal (ex.: Iterative Cone)

- iterate from the hardest seed
- remove the stable cone as a jet and start again

Idea: “regular/circular” jets

## Recombination:

- $k_t$  algorithm
- Cambridge/Aachen alg.

## Cone:

- CDF JetClu
- CDF MidPoint
- D0 (run II) Cone
- PxCone
- ATLAS Cone
- CMS Iterative Cone
- PyCell/CellJet
- GetJet

***Part 0***  
***21st century: towards a solid toolkit***

SNOWMASS accords, Tevatron 1990 (i.e. old!):

Several important properties that should be met by a jet definition are [3]:

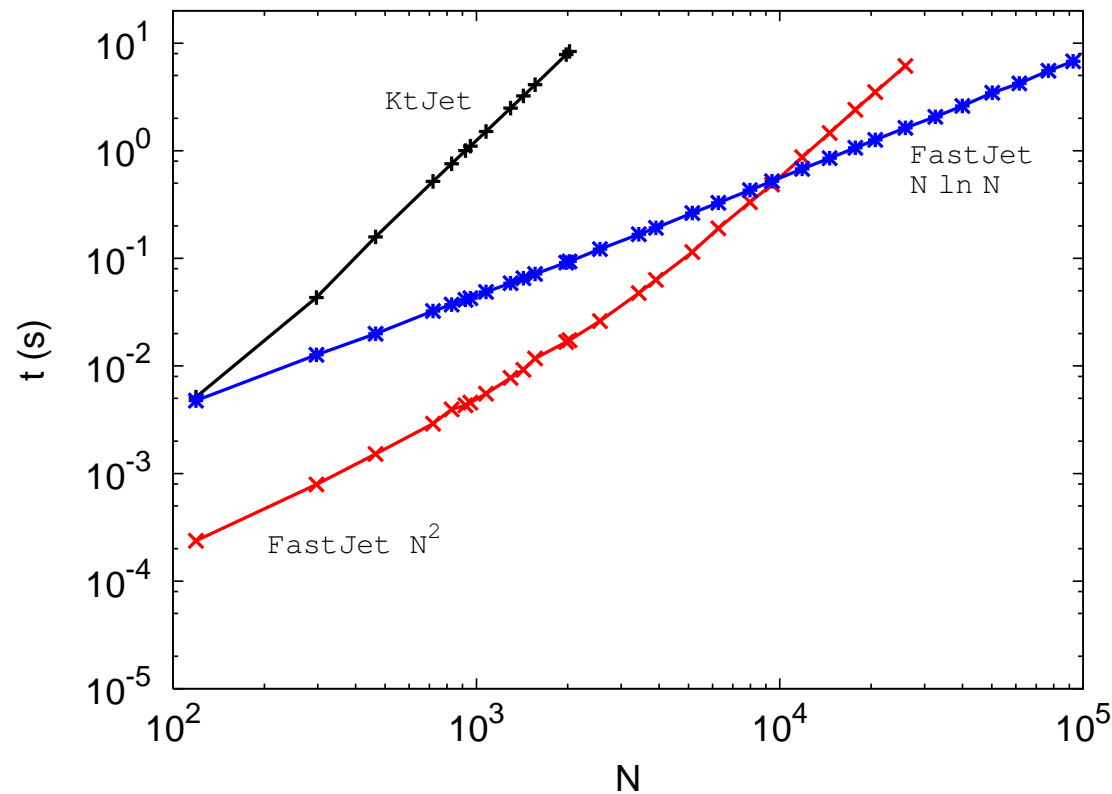
1. Simple to implement in an experimental analysis;
2. Simple to implement in the theoretical calculation;
3. Defined at any order of perturbation theory;
4. Yields finite cross section at any order of perturbation theory;
5. Yields a cross section that is relatively insensitive to hadronization.

i.e. usable by theoreticians (e.g. finite perturbative results)  
and experimentalists (e.g. fast enough, not much UE sensitivity)

[M. Cacciari, G. Salam, 06]

## Speeding up the $k_t$ and Cam/Aachen algorithms

- using computational-geometry techniques:  $\mathcal{O}(N^3) \rightarrow \mathcal{O}(N \log N)$
- C++ implementation in FastJet



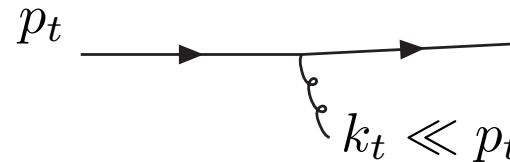
QCD probability for gluon bremsstrahlung at angle  $\theta$  and  $\perp$ -mom.  $k_t$ :

$$dP \propto \alpha_s \frac{d\theta}{\theta} \frac{dk_t}{k_t}$$

Two divergences:



Collinear



Soft



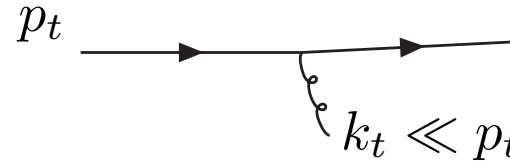
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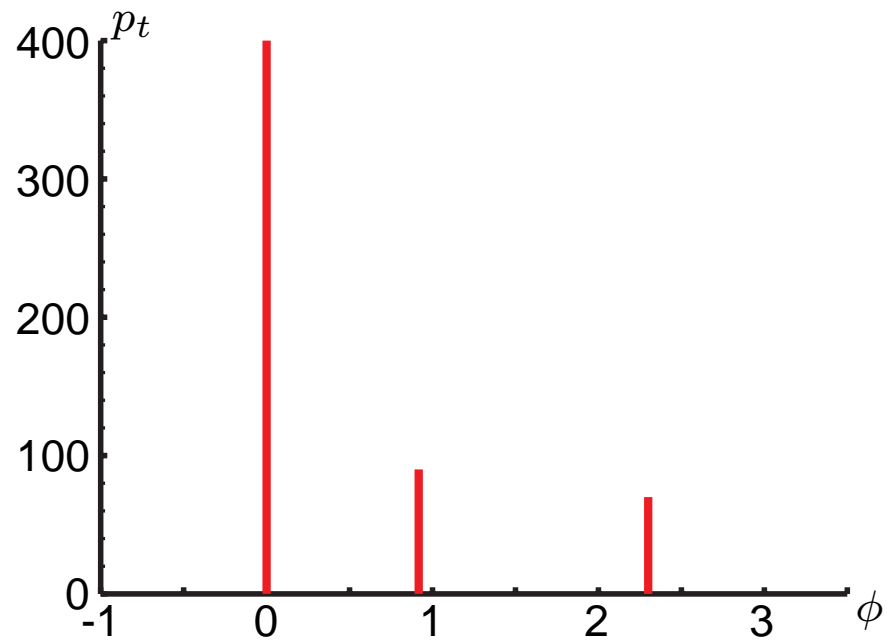
Collinear

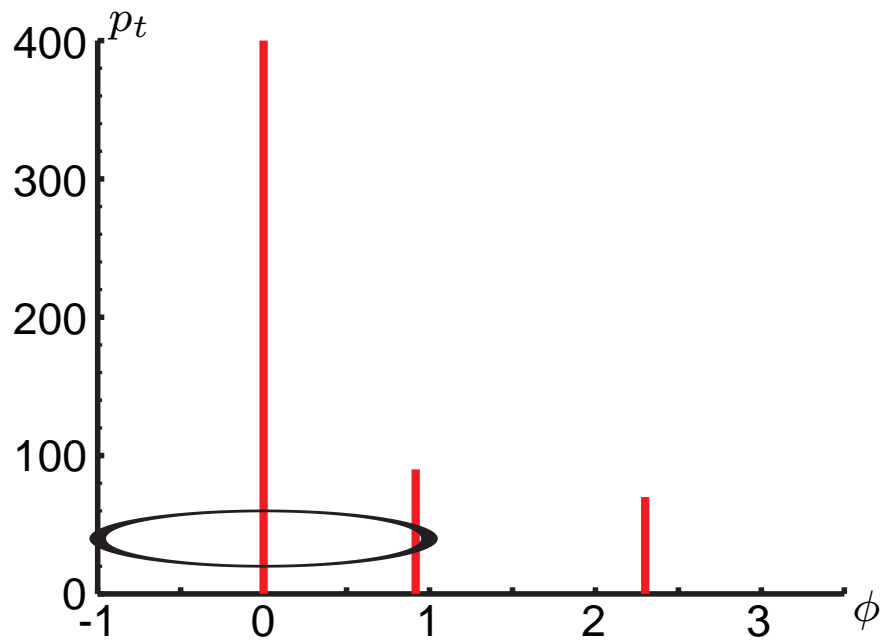


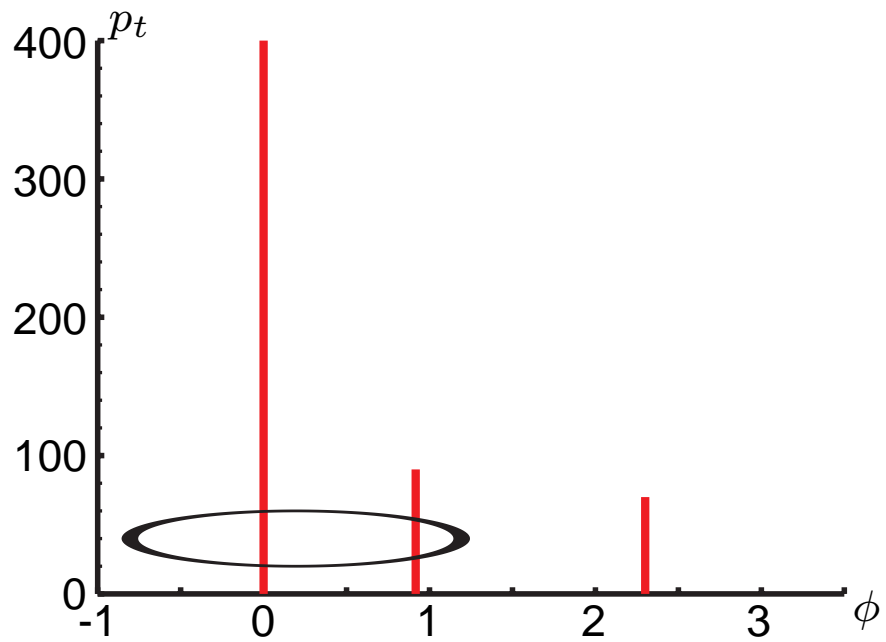
Soft

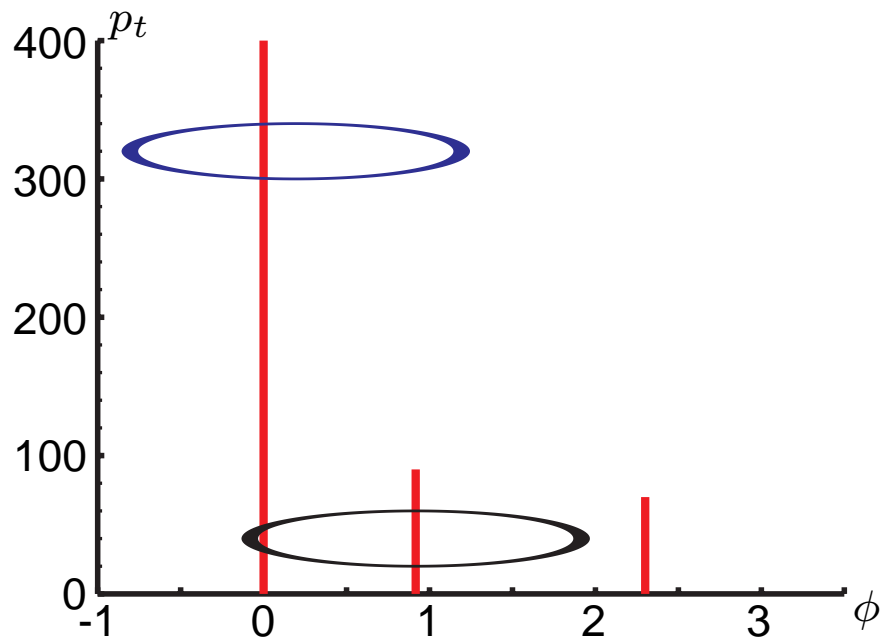
For pQCD to make sense, the (hard) jets should not change when

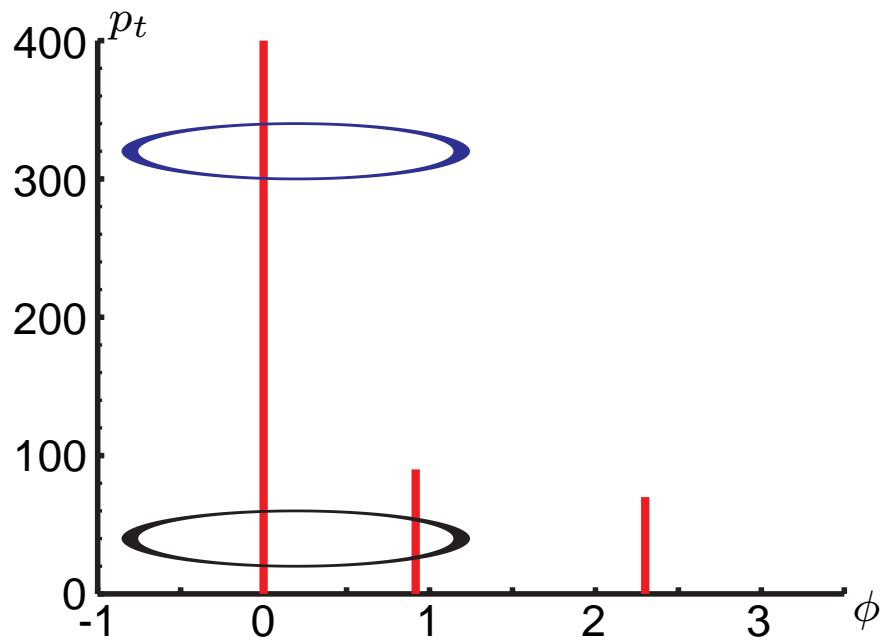
- one has a collinear splitting  
*i.e.* replaces one parton by two at the same place  $(\eta, \phi)$
- one has a soft emission *i.e.* adds a very soft gluon

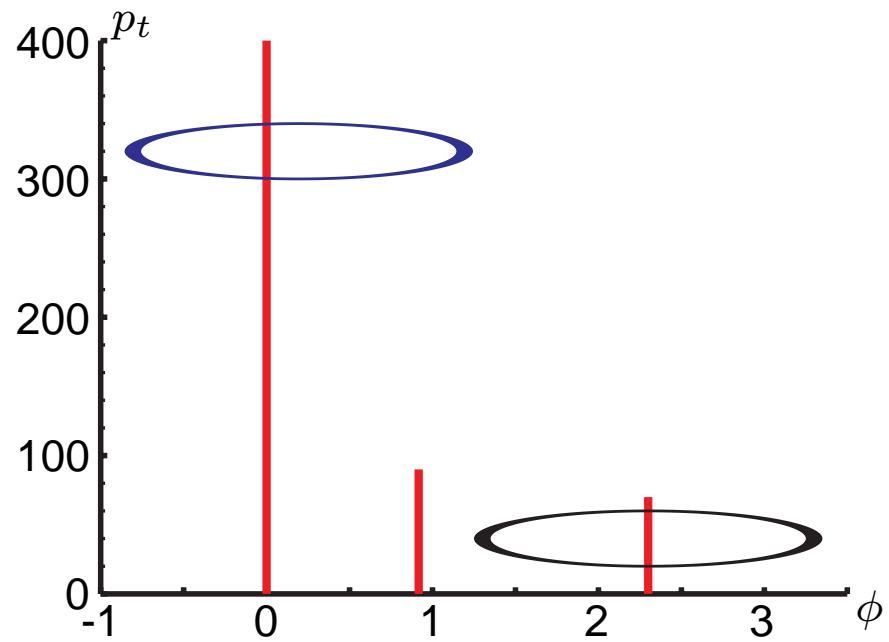


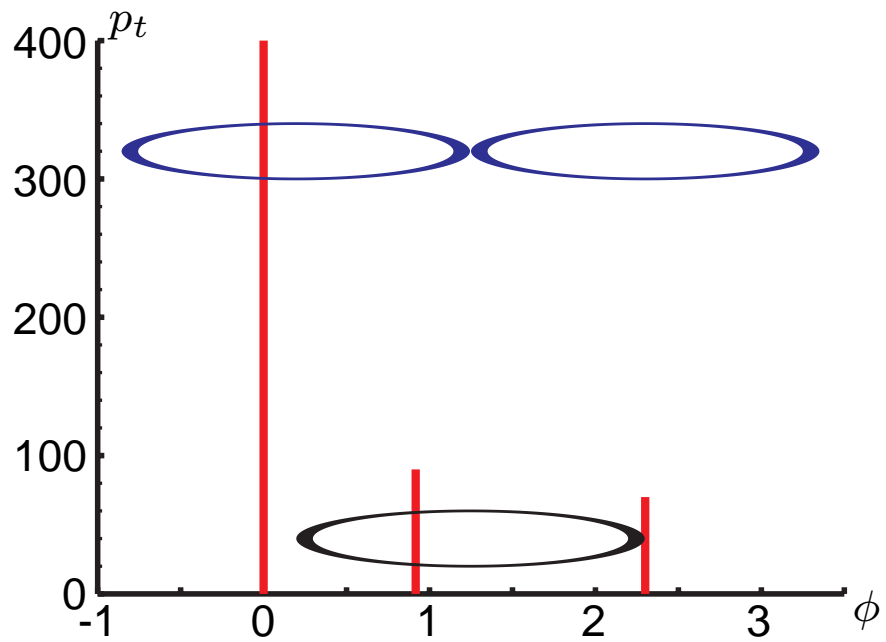






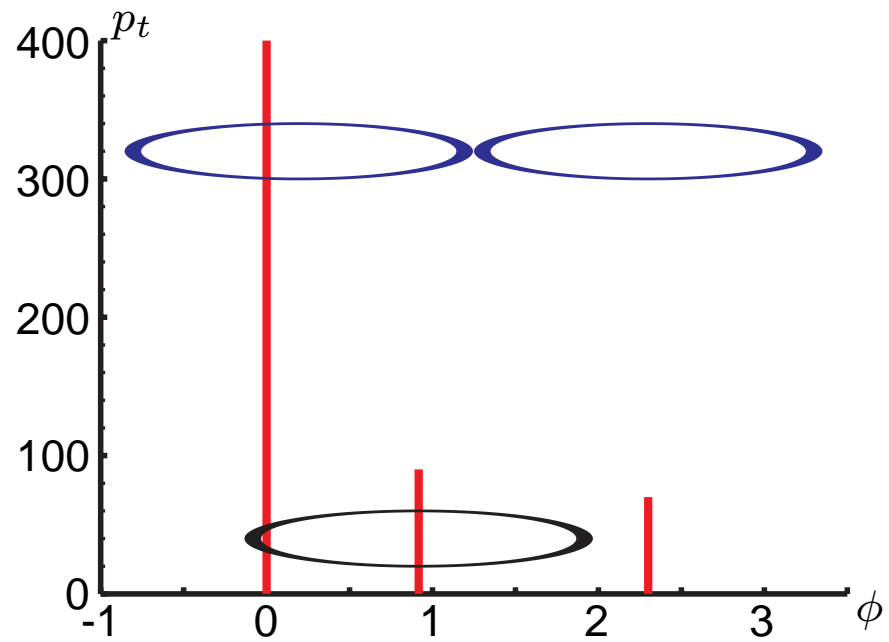




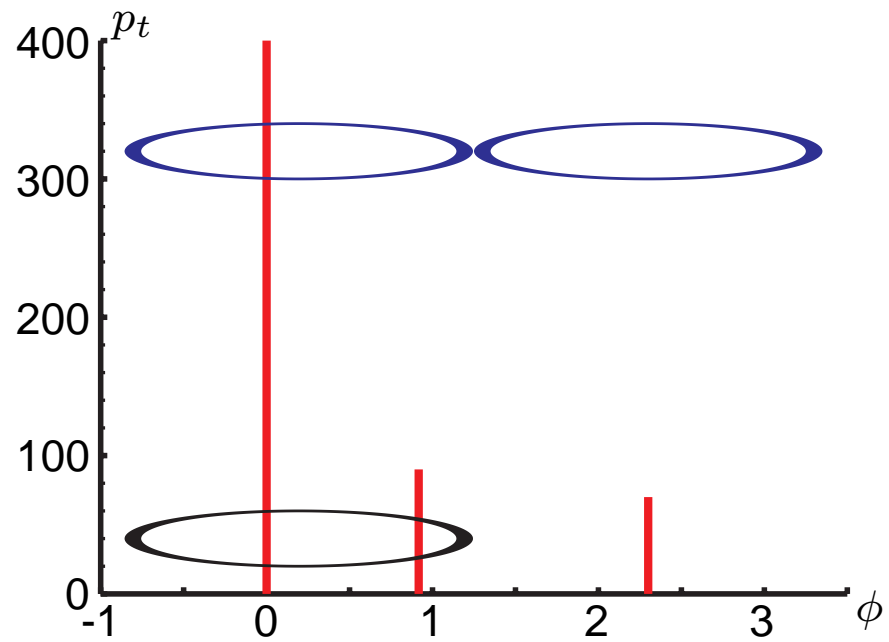




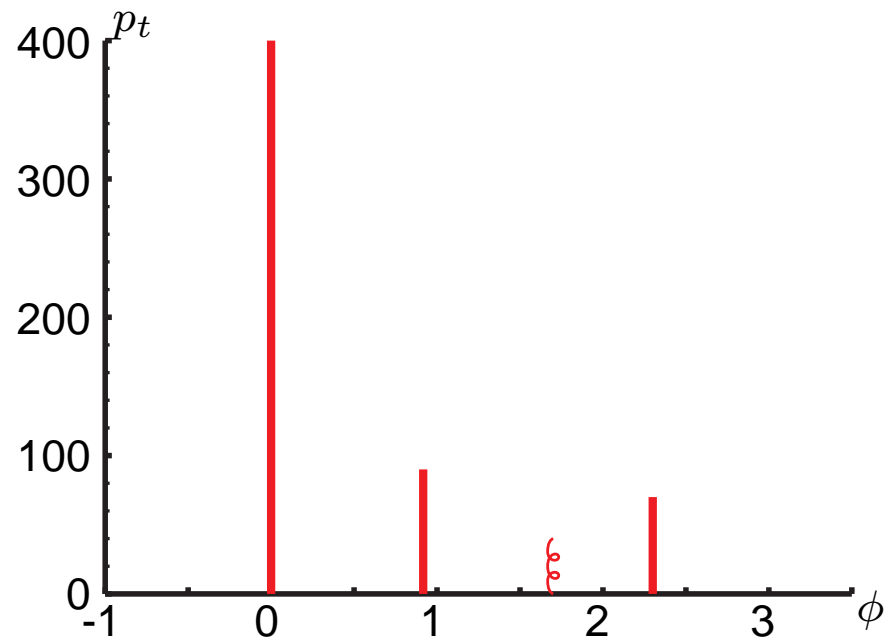
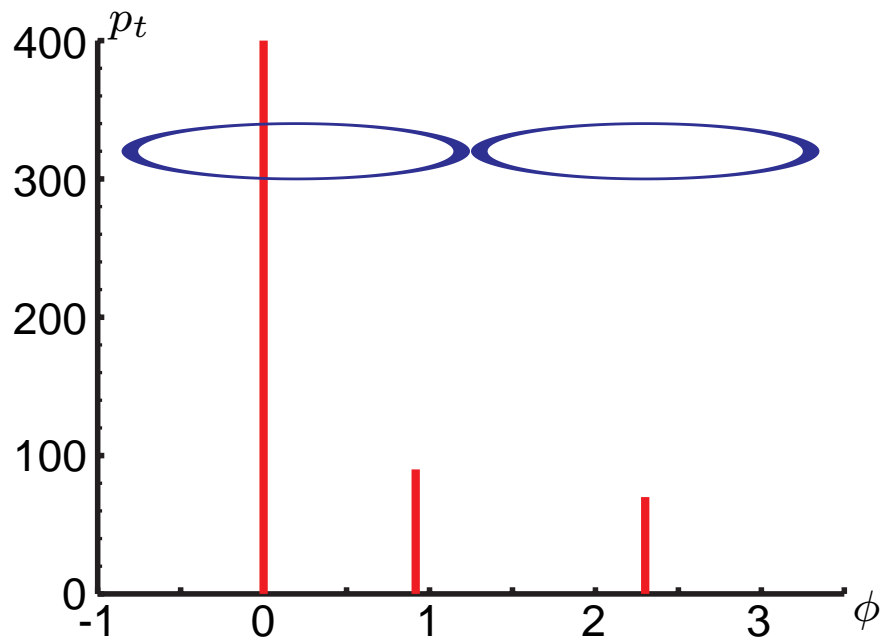
# IR unsafety of the Midpoint alg



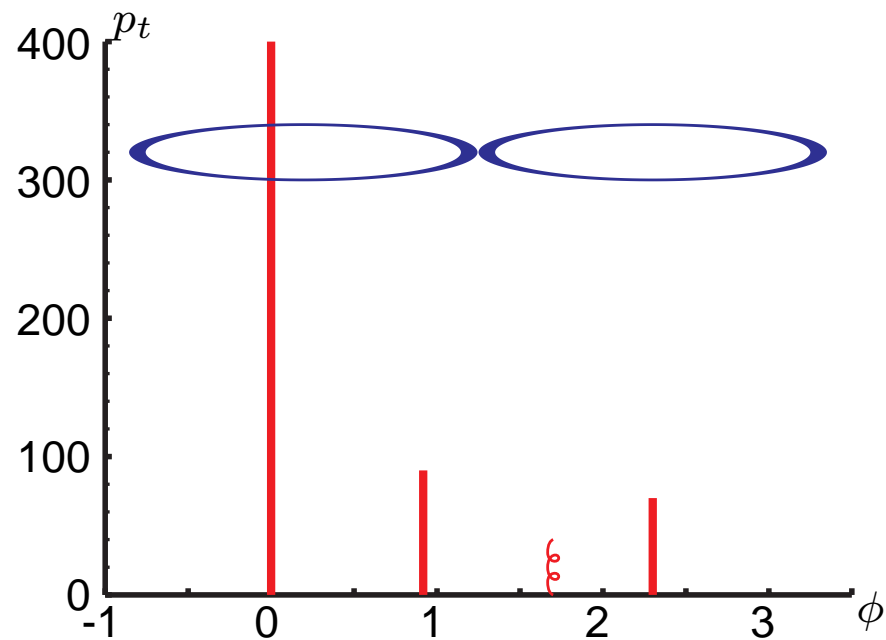
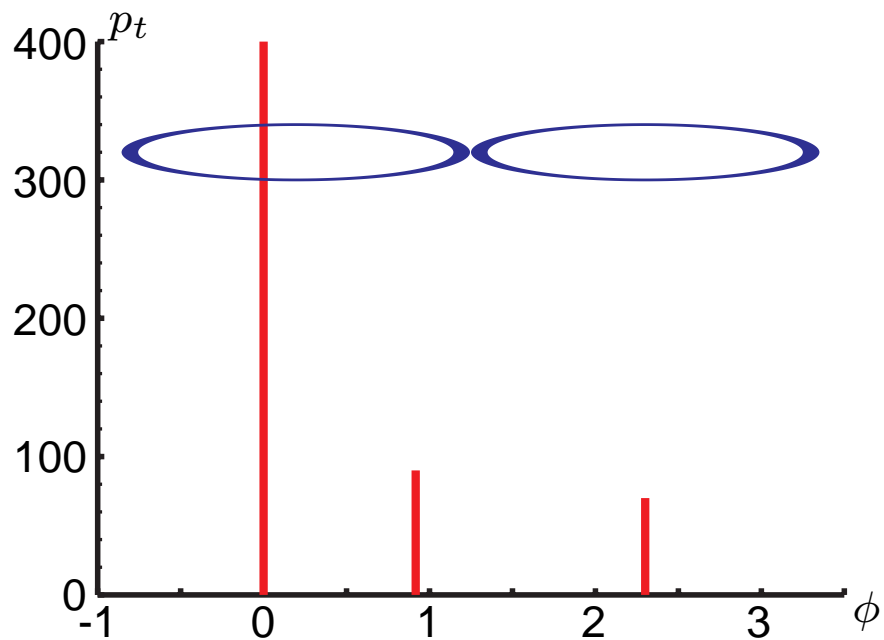
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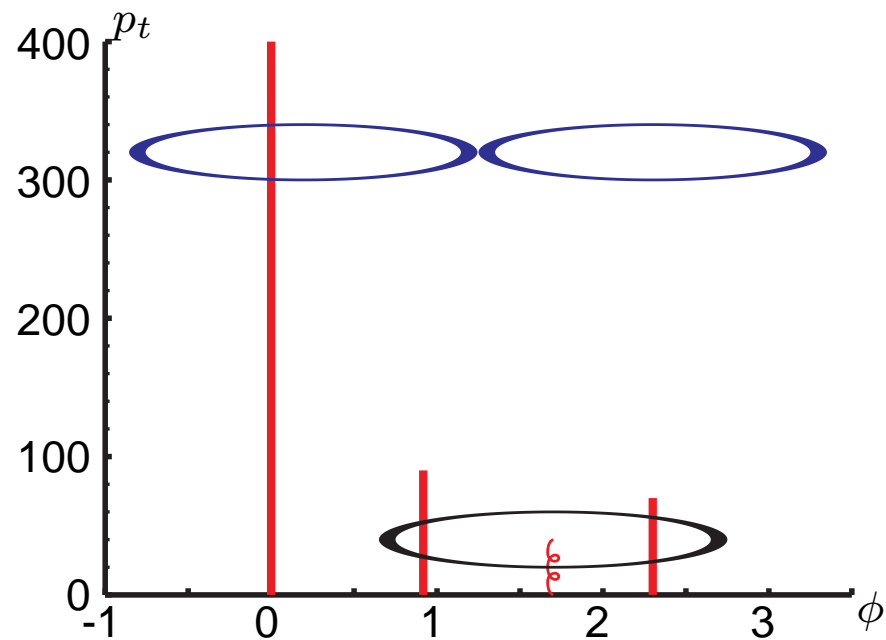
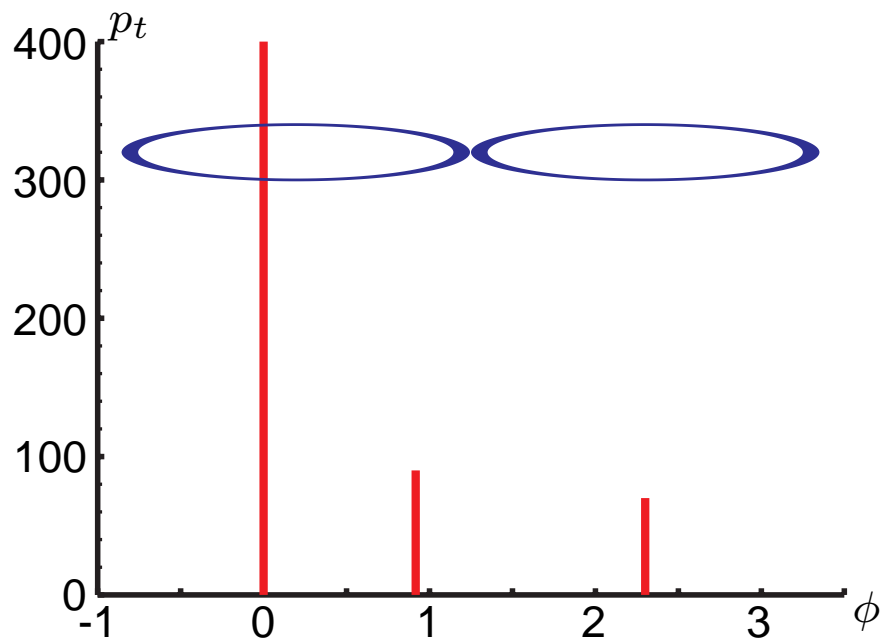
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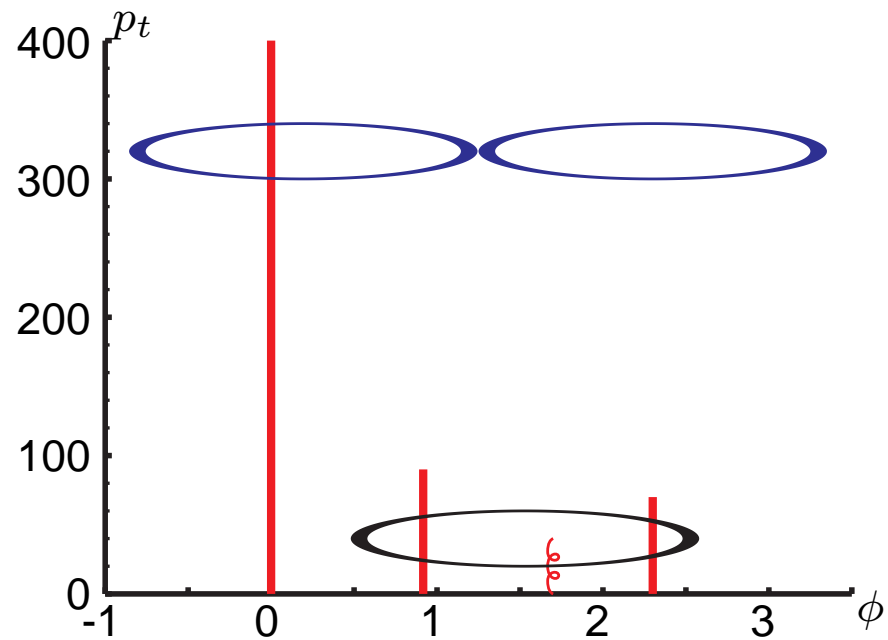
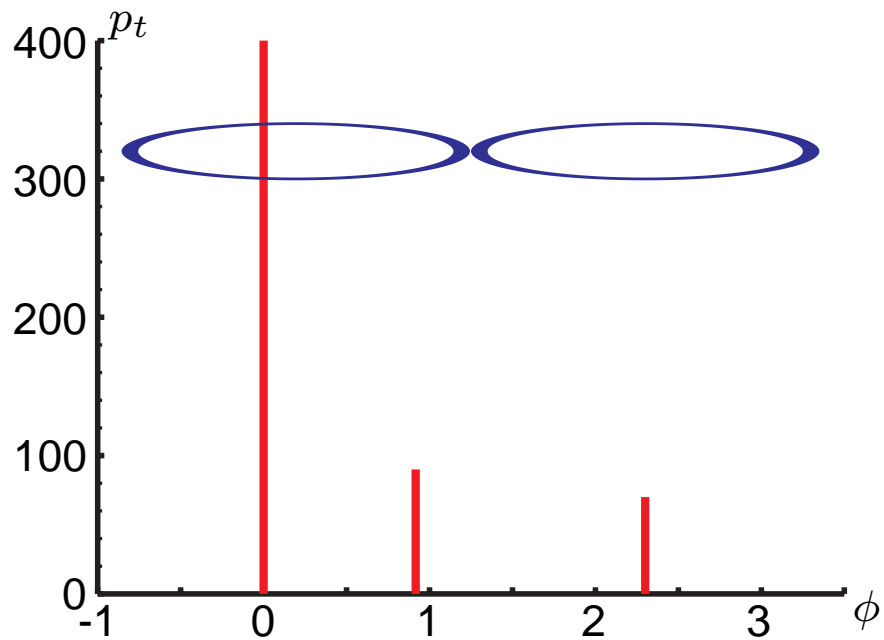
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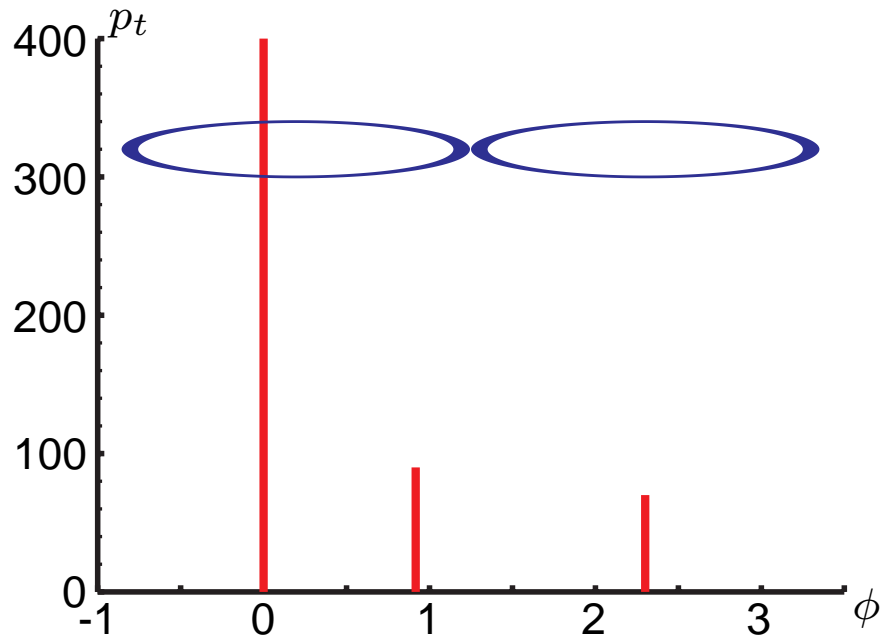


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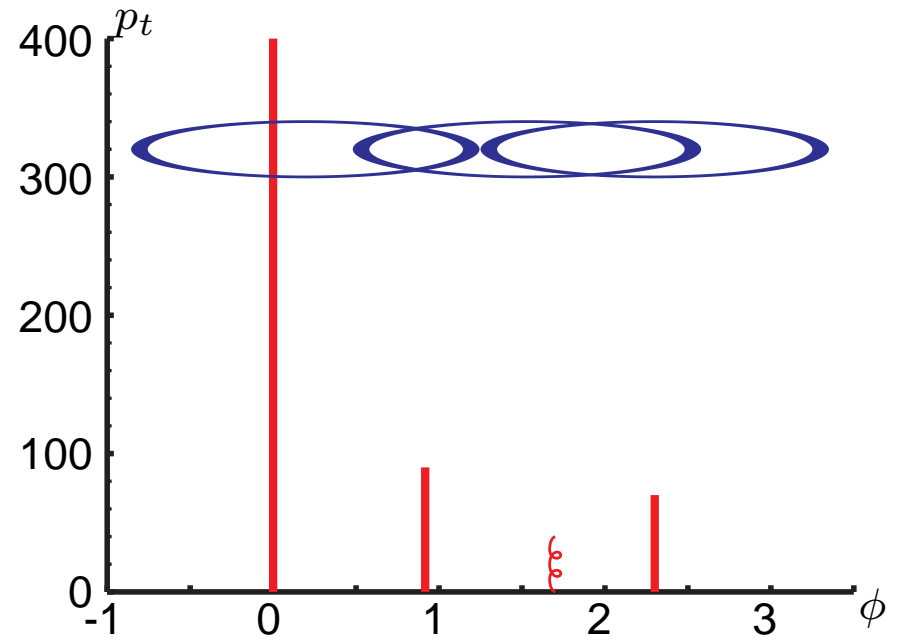
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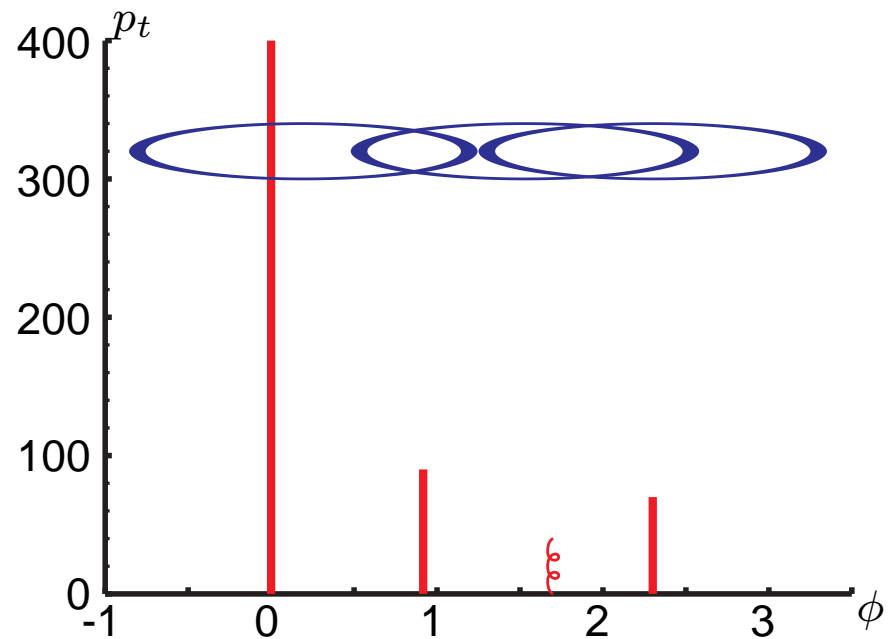
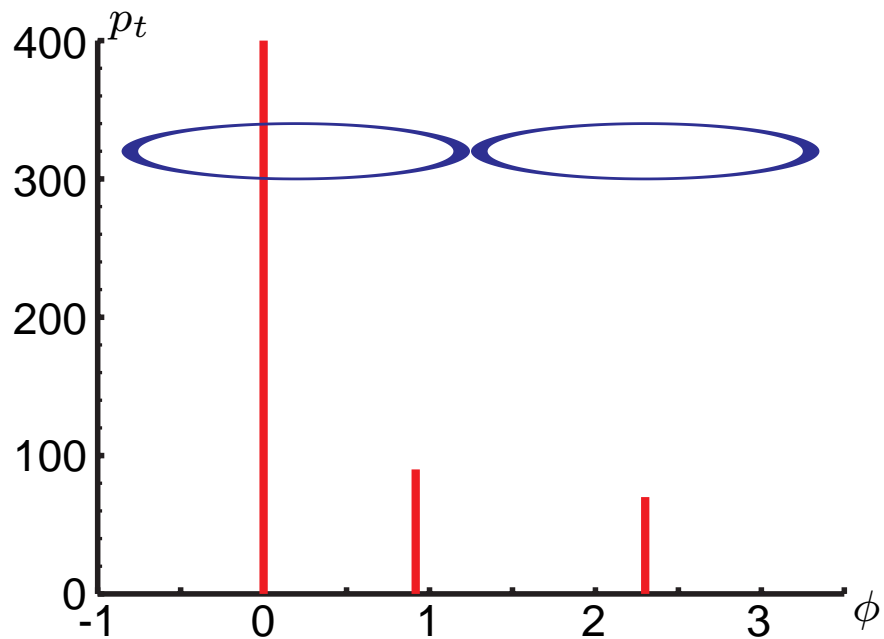


Stable cones:

Midpoint:                    {1,2} & {3}



{1,2} & {3} & {2,3}



Stable cones:

Midpoint:  $\{1,2\}$  &  $\{3\}$

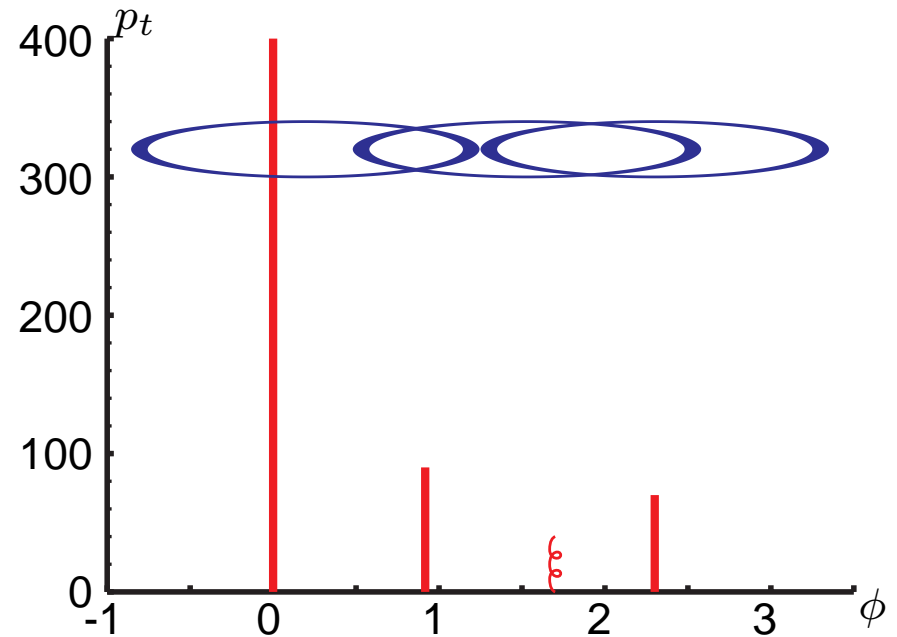
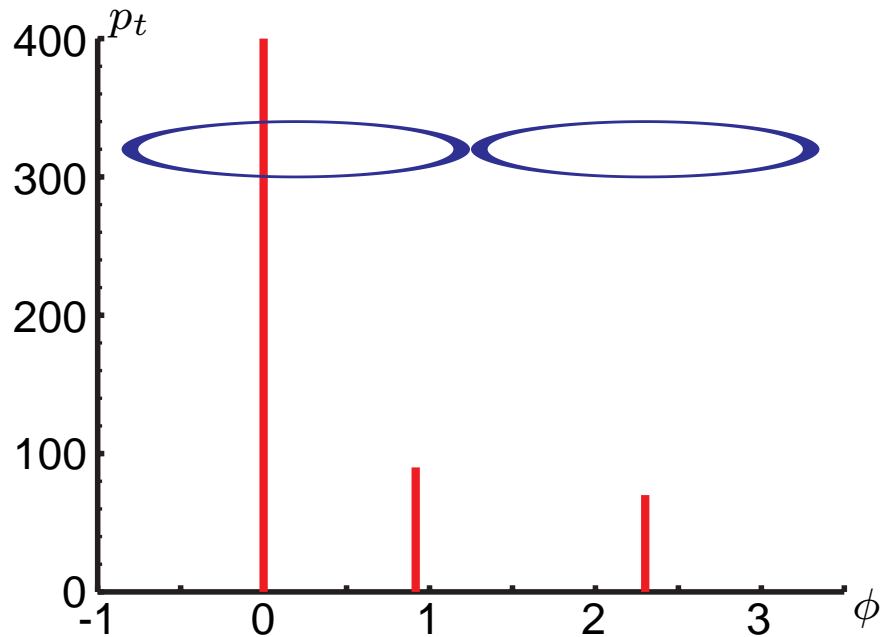
$\{1,2\}$  &  $\{3\}$  &  $\{2,3\}$

Jets: ( $f = 0.5$ )

Midpoint:  $\{1,2\}$  &  $\{3\}$

$\{1,2,3\}$





Stable cones:

Midpoint: {1,2} & {3}

Seedless: {1,2} & {3} & {2,3}

{1,2} & {3} & {2,3}

{1,2} & {3} & {2,3}

Jets: ( $f = 0.5$ )

Midpoint: {1,2} & {3}

Seedless: {1,2,3}

{1,2,3}

{1,2,3}

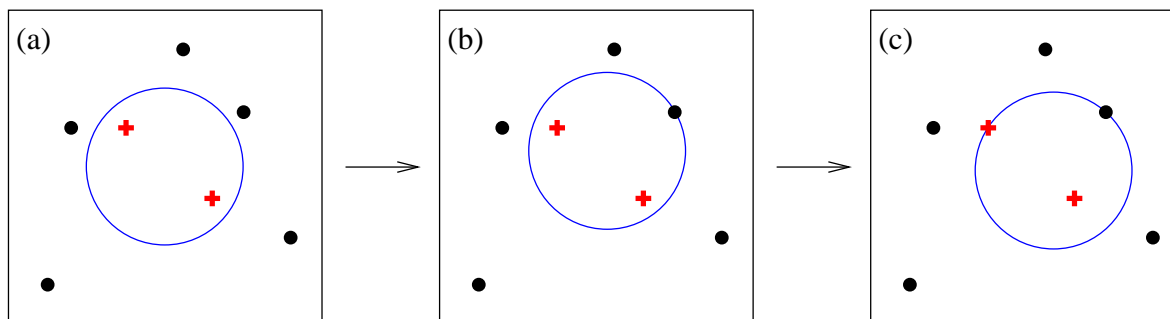
**Stable cone missed → IR unsafety of the midpoint algorithm**

- Solution: use a seedless approach, find **ALL** stable cones
- Naive approach: check stability of each subset of particle

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- Naive approach: check stability of each subset of particle  
Complexity is  $\mathcal{O}(N2^N)$   
 $\Rightarrow$  **definitely unrealistic:  $10^{17}$  years for  $N = 100$**
- Midpoint complexity:  $\mathcal{O}(N^3)$

- Solution: use a seedless approach, find **ALL** stable cones
- Midpoint complexity:  $\mathcal{O}(N^3)$

Idea: use geometric arguments



- Each enclosure can be moved (in any direction) until it touches a point
- ... then rotated until it touches a second one

- Solution: use a seedless approach, find **ALL** stable cones
- Midpoint complexity:  $\mathcal{O}(N^3)$

Idea: use geometric arguments

⇒ Enumerate all pairs of particles  
with 2 circle orientations and 4 possible inclusion/exclusion  
→ find all enclosures

- Complexity:  $\mathcal{O}(N^3)$ , with improvements:  $\mathcal{O}(N^2 \log(N))$

- Solution: use a seedless approach, find **ALL** stable cones
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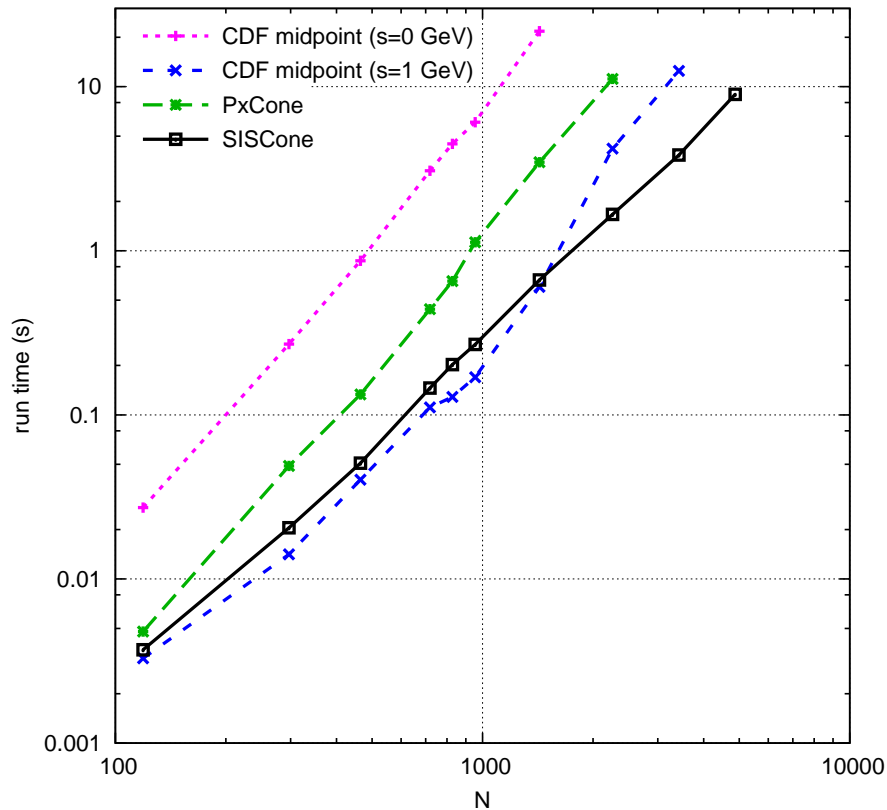
- Complexity:  $\mathcal{O}(N^3)$ , with improvements:  $\mathcal{O}(N^2 \log(N))$

→ C++ implementation: Seedless Infrared-Safe Cone algorithm (SIScone)  
G.Salam, G.S., JHEP 04 (2007) 086; <http://projects.hepforge.org/siscone>

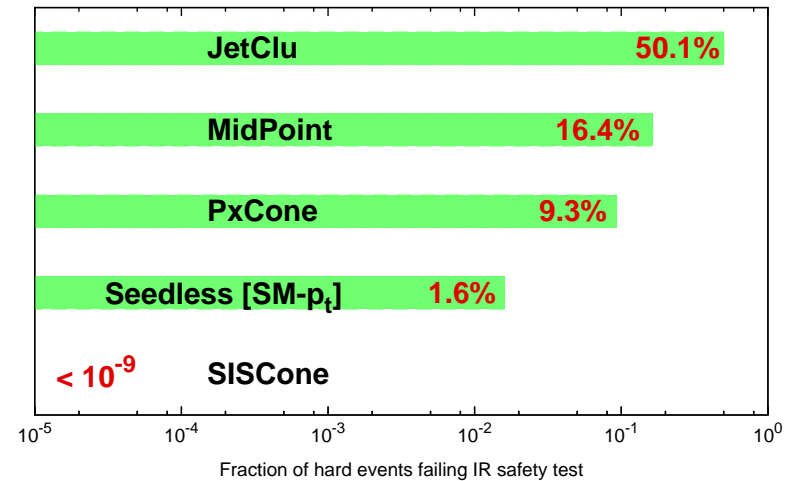
NB.: also available from FastJet

[M.Cacciari, G.Salam, G.S.]; <http://www.lpthe.jussieu.fr/~salam/fastjet>

## Execution timings:



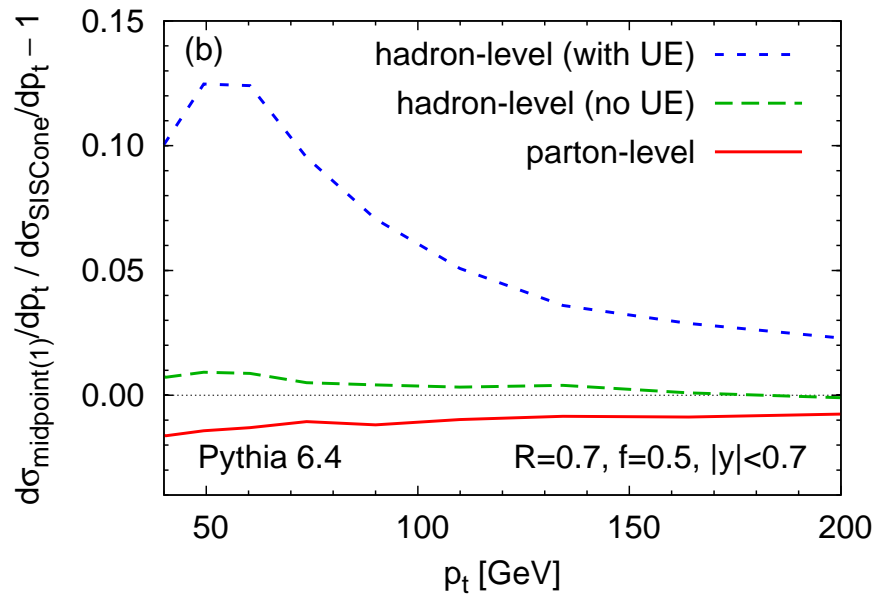
Random hard & soft partons  
fraction with “hard  $\neq$  hard+soft”



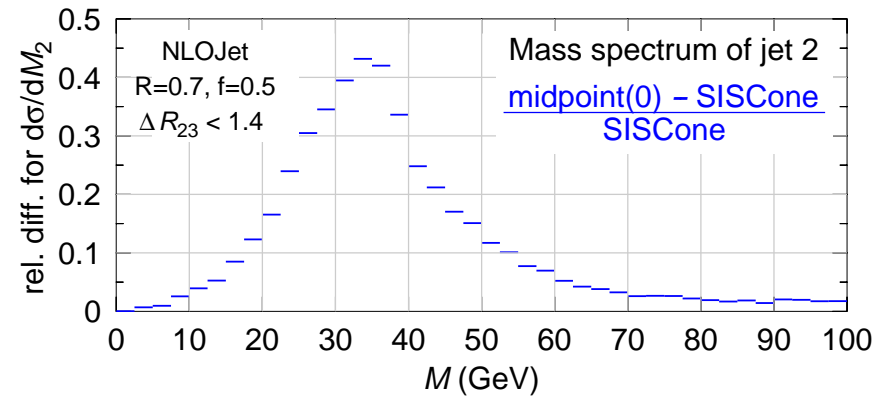
- at least as fast as midpoint cones
- IR safe
  - JetClu, ATLAS cone: 50% failure
  - MidPoint: 15% failure

## Inclusive (midpoint/SISCone-1)

pp  $\sqrt{s} = 14$  TeV



## Masses in 3-jet events

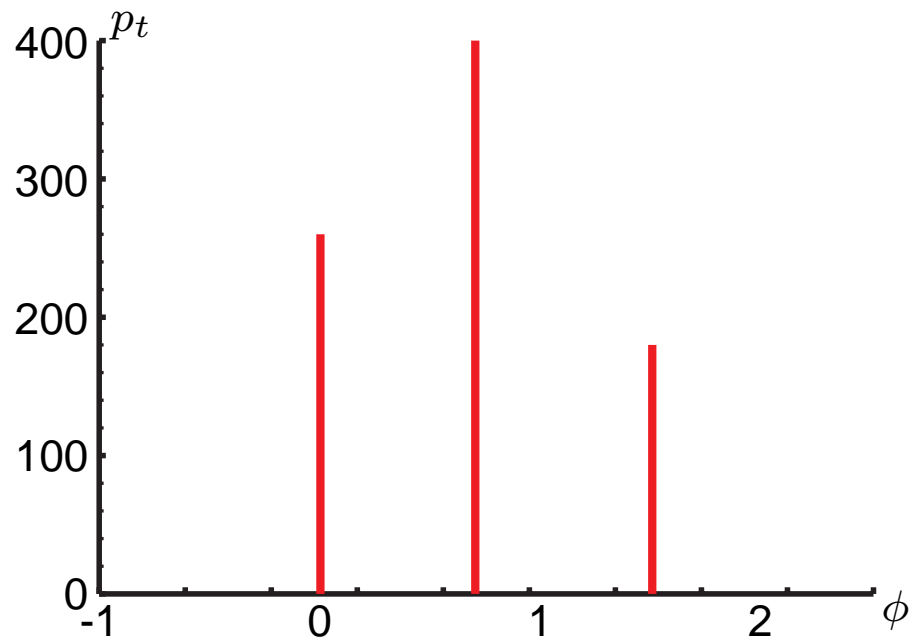


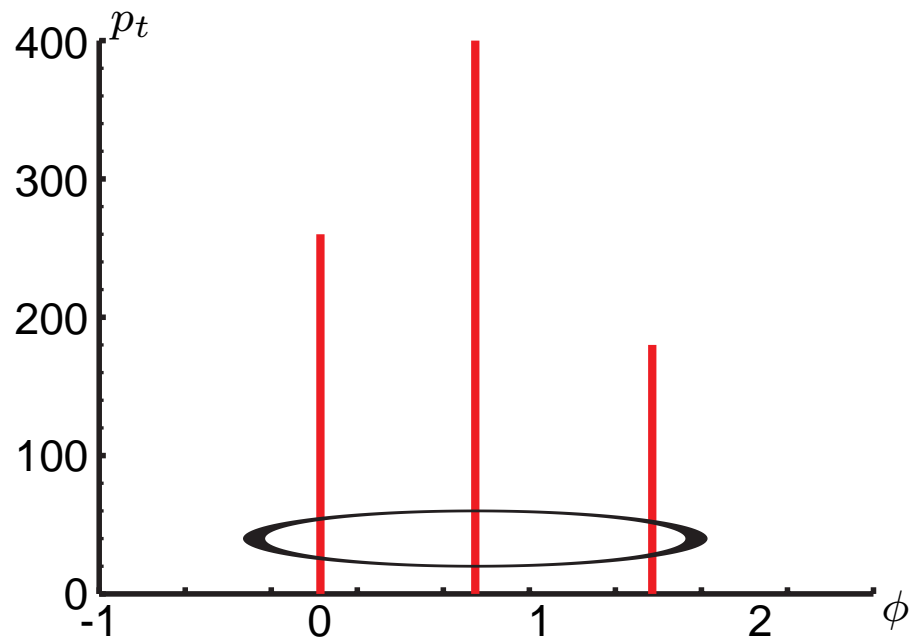
### Inclusive cross-section:

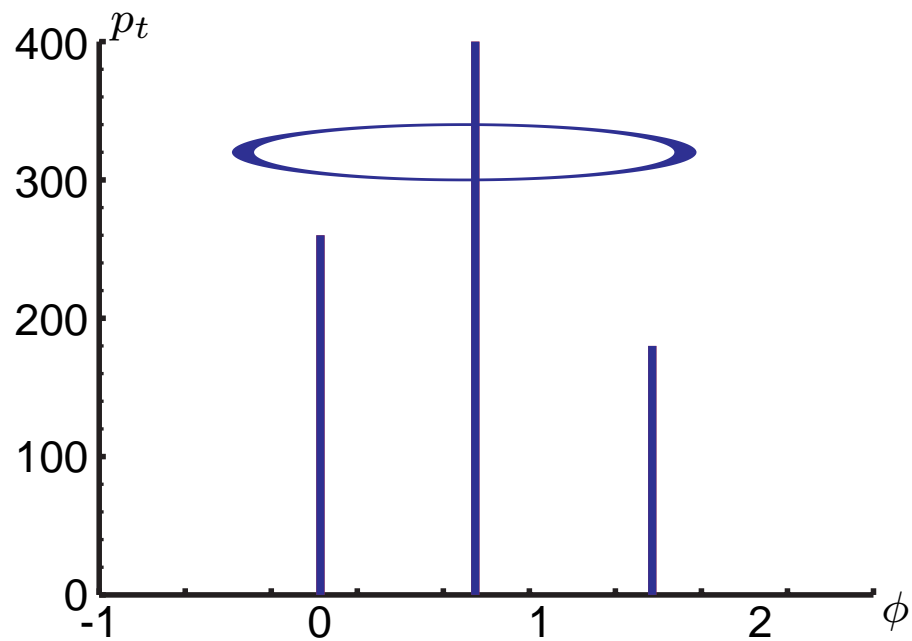
- Effect of a few percents
- Less sensitivity to the UE

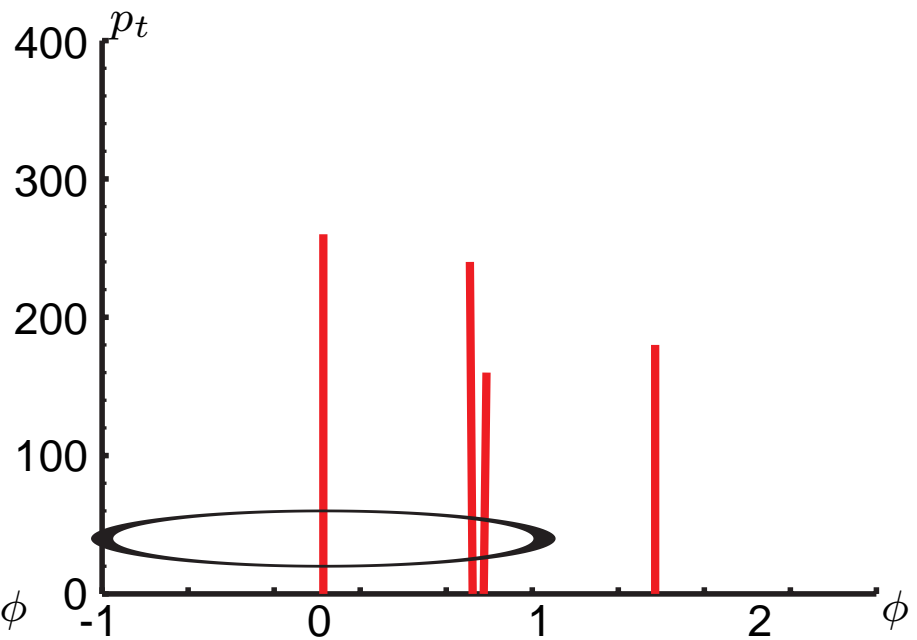
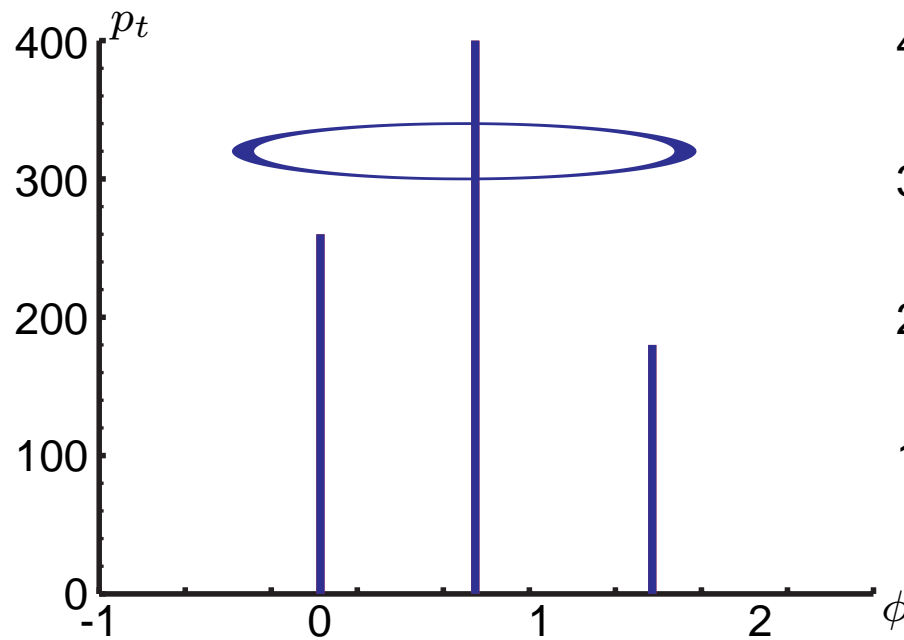
• More exclusive processes: effects  $\sim 45\%$  (Important for LHC!)

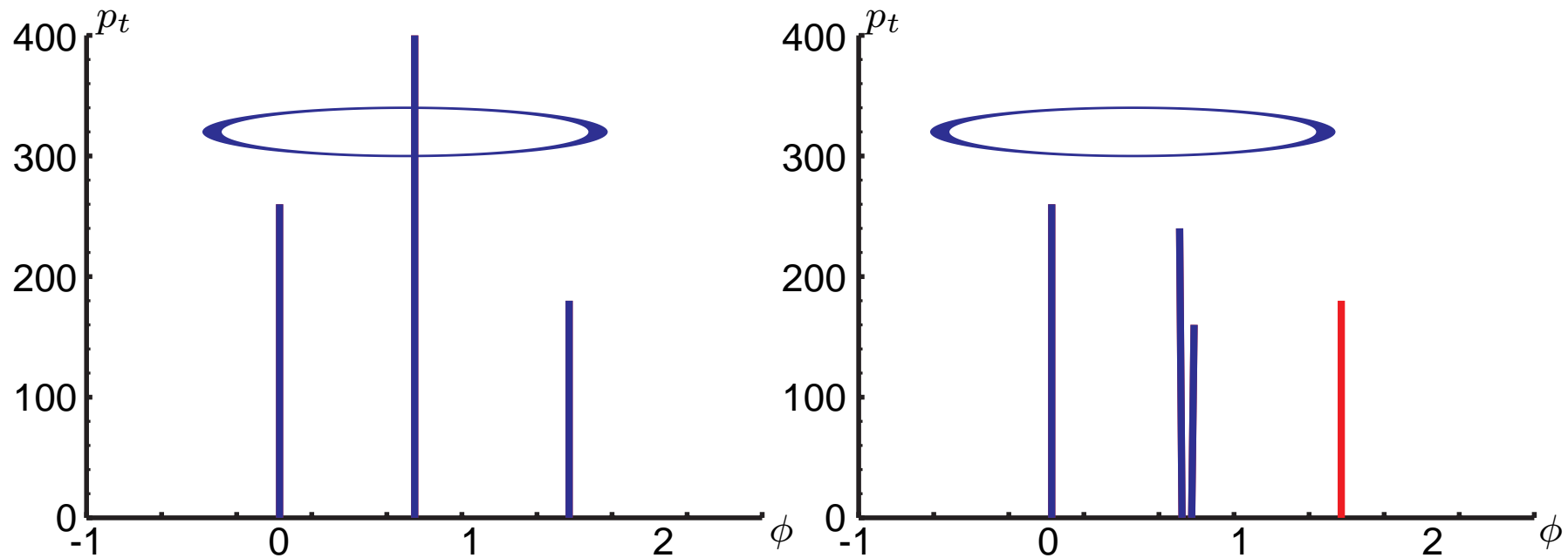


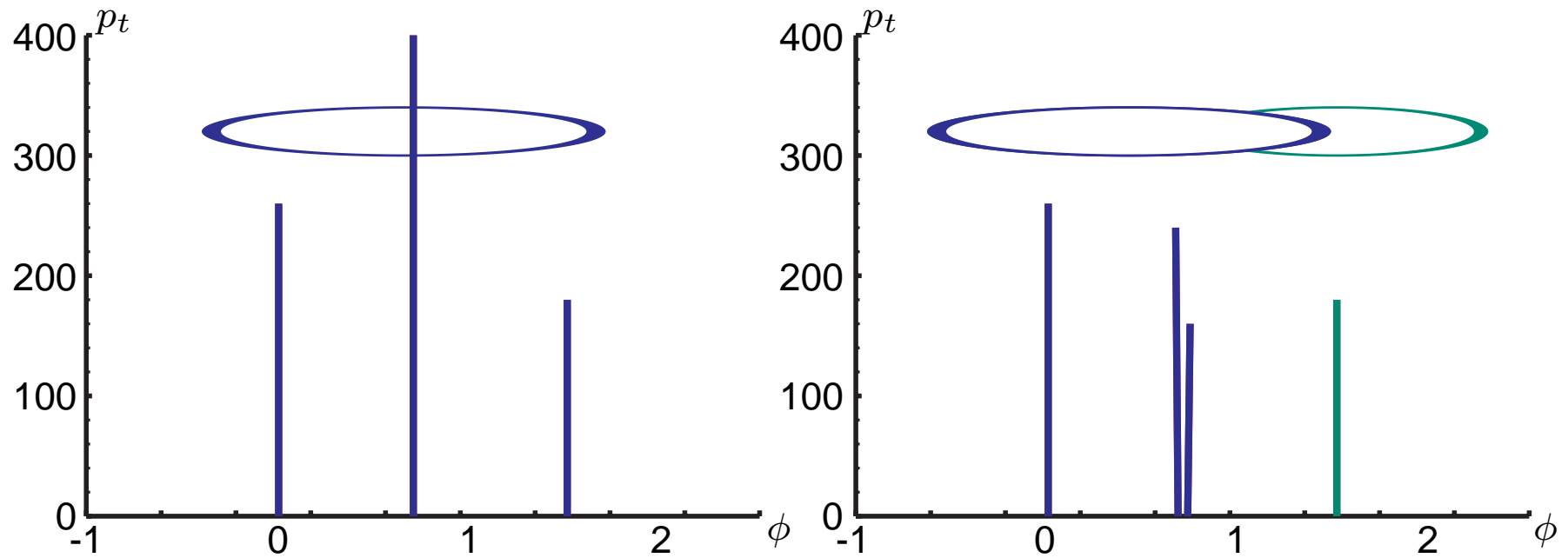


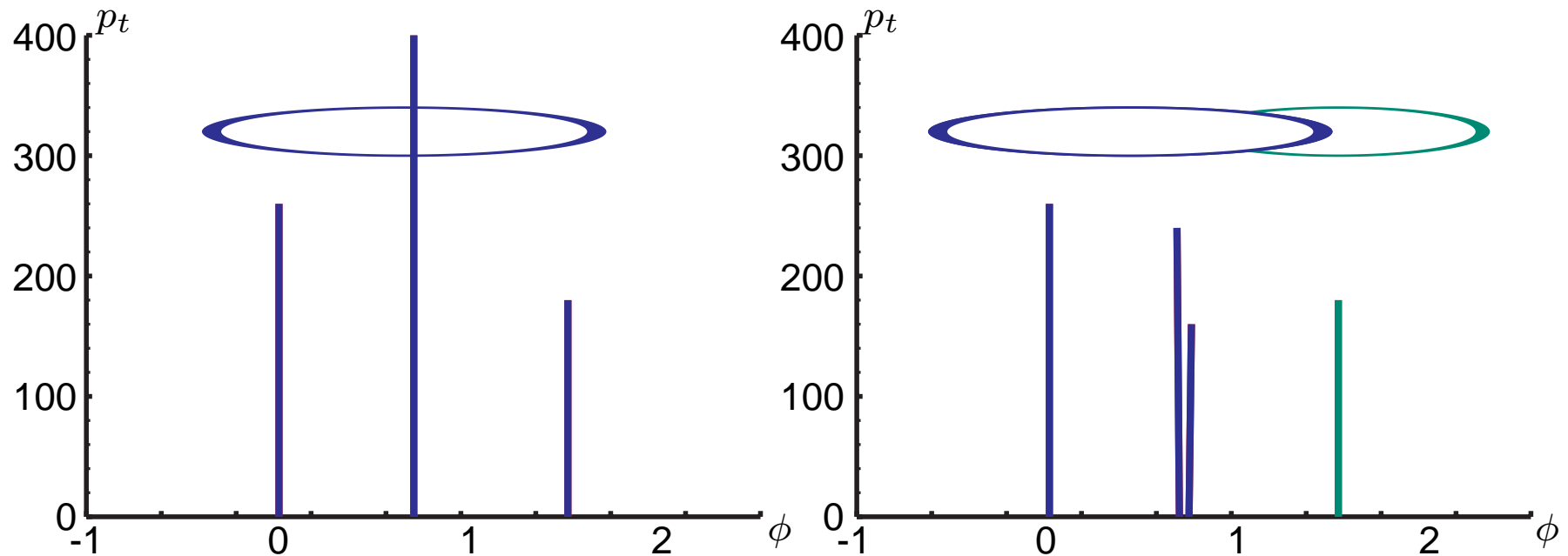












- Before collinear splitting: 1 jet
- After collinear splitting: 2 jets

→ **collinear unsafety of the iterative cone algorithm**

Come back to recombination-type algorithms:

$$d_{ij} = \min(k_{t,i}^{2p}, k_{t,j}^{2p}) (\Delta\phi_{ij}^2 + \Delta\eta_{ij}^2)$$

- $p = 1$ :  $k_t$  algorithm
- $p = 0$ : Aachen/Cambridge algorithm



Come back to recombination-type algorithms:

$$d_{ij} = \min(k_{t,i}^{2p}, k_{t,j}^{2p}) (\Delta\phi_{ij}^2 + \Delta\eta_{ij}^2)$$

- $p = 1$ :  $k_t$  algorithm
- $p = 0$ : Aachen/Cambridge algorithm
- $p = -1$ : anti- $k_t$  algorithm [M.Cacciari, G.Salam, G.S., JHEP 04 (08) 063]

Come back to recombination-type algorithms:

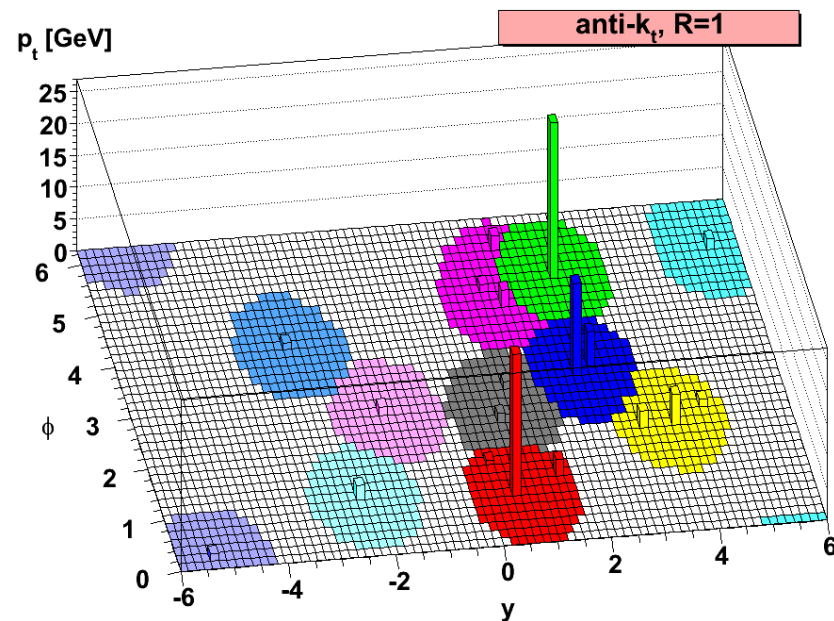
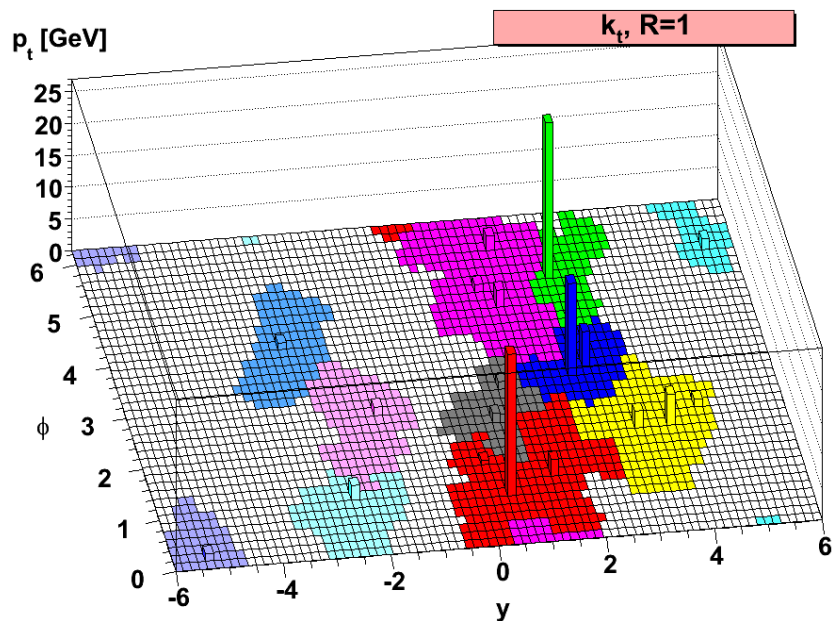
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Why should that be related to the iterative cone ?!?

- “large  $k_t \Rightarrow$  small distance”  
*i.e.* hard partons “eat” everything up to a distance  $R$   
*i.e.* circular/regular jets, jet borders unmodified by soft radiation
- infrared and collinear safe

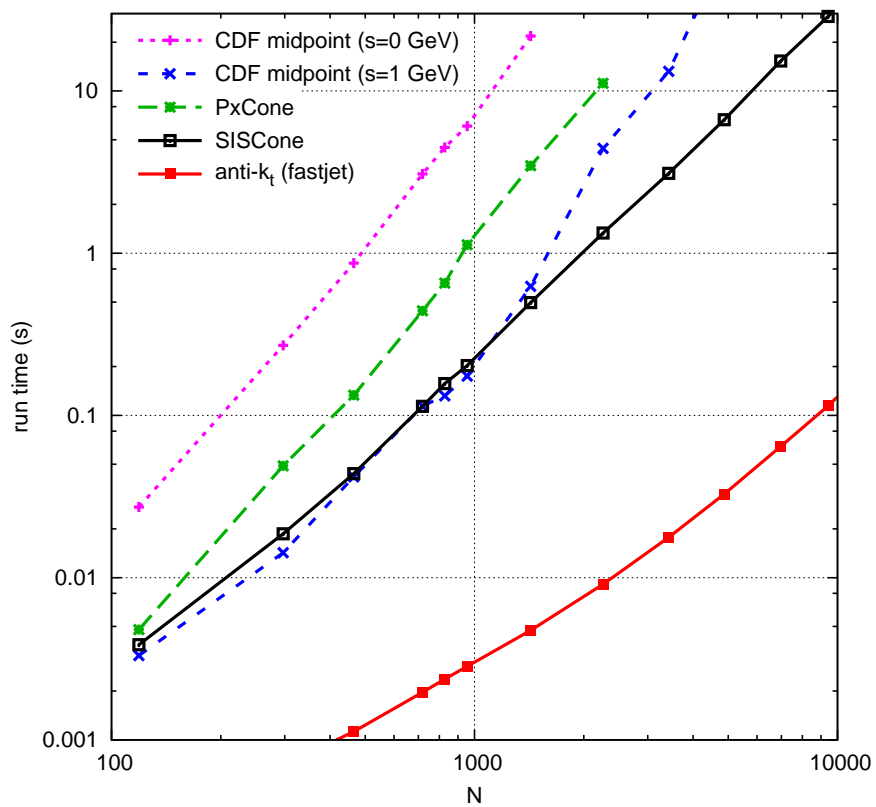
## Hard event + homogeneous soft background



anti- $k_t$  is soft-resilient

more later in this talk...

## Execution timings:



As fast as the (fast)  $k_t$  ([M. Cacciari, G. Salam, 06])

Midpoint and the iterative cone IR or Collinear unsafe (at  $\mathcal{O}(\alpha_s^4)$ )

Observable	1st miss cones at	Last meaningful order
Inclusive jet cross section	NNLO	NLO
3 jet cross section	NLO	LO (NLO in NLOJet)
$W/Z/H + 2$ jet cross sect.	NLO	LO (NLO in MCFM)
jet masses in 3 jets	LO	none (LO in NLOJet)

⇒ The IRC-unsafety issue will matter at LHC

+ We do not want the theoretical efforts to be wasted

Note: 1 order worse for JetClu of the ATLAS Cone!

## Recombination:

- $k_t$  algorithm
- Cambridge/Aachen alg.
- **anti- $k_t$  algorithm**

4 available  
safe algorithms

All part of FastJet

[M.Cacciari,G.Salam,G.S.]

## Cone:

- CDF JetClu
- CDF MidPoint
- D0 (run II) Cone
- PxCone
- ATLAS Cone
- CMS Iterative Cone
- PyCell/CellJet
- GetJet
- **SISCone**

More refined clustering (“2<sup>nd</sup> generation of algorithms”)

Cambridge+Filtering algorithm:

- Cluster with Aachen/Cambridge and radius  $R$
- For each jet, recluster it with Aachen/Cambridge and radius  $R_{\text{sub}}$   
keep only  $n_{\text{sub}}$  hardest sub-jets of the initial jet

Aim: remove the soft background

Properties:

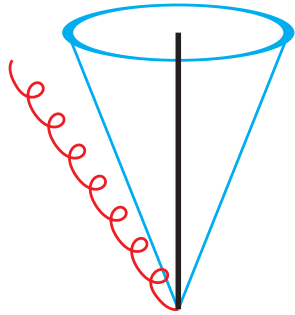
- Proven to improve jet reconstruction, in  $H \rightarrow b\bar{b}$   
[J.Butterworth, A.Davison, M.Rubin, G.Salam, 08]
- Additional parameters that deserve appropriate studies
- We will use the simplest choice:  $R_{\text{sub}} = R/2$ ,  $n_{\text{sub}} = 2$

***Part 1***  
***Jets in  $pp$  collisions***  
***Choosing the adapted jet definition***



Competition between

- catching perturbative radiation



Out-of-jet radiation:

$$\sim \int_R \frac{d\theta}{\theta} \sim \log(1/R)$$

- not catching soft background radiation

$$\text{Soft contents} \sim \text{jet area} \sim R^2$$

more detailed analytic computation in progress...

only numerical results at present time

We analyse 3 processes:

- $Z' \rightarrow q\bar{q} \rightarrow 2 \text{ jets}$ : (fictitious narrow  $Z'$ )  
simple environment: identify 2 jets and reconstruct  $M_{Z'}$   
source of monochromatic quark jets  
scale dependence: mass of the  $Z'$  between 100 GeV and 4 TeV
- $H \rightarrow gg \rightarrow 2 \text{ jets}$ : (fictitious narrow Higgs)  
simple environment: identify 2 jets and reconstruct  $M_H$   
source of monochromatic gluon jets  
scale dependence: mass of the Higgs between 100 GeV and 4 TeV
- $t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow q\bar{q}bq\bar{q}\bar{b} \rightarrow 6 \text{ jets}$ :  
complex environment: identify 6 jets and reconstruct 2 top  
balance between reconstruction efficiency and identification

with

- the 5 IRC-safe algorithms:  $k_t$ , Cambridge, anti- $k_t$ , SISCone, Cam+filtering
- jet radius varied between 0.1 and 1.5

## Measure of the jet reconstruction efficiency

- Forget about measures related to parton-jet matching, use the reconstructed mass peak
- Forget about fits depending on the shape of the peak
- we shall maximise the signal over background ratio ( $S/\sqrt{B}$ ):

$Q_{f=z}^w(JA, R)$  = minimal width of a window containing a fraction  $f = z$  of the events

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- it intuitively does what it should
- for a constant background,

$$\frac{Q_{f=z}^w(JA_1, R_1)}{Q_{f=z}^w(JA_2, R_2)} = \frac{B_{JA_1, R_1}}{B_{JA_2, R_2}} = \left[ \frac{(S/\sqrt{B})_{JA_1, R_1}}{(S/\sqrt{B})_{JA_2, R_2}} \right]^{-2}$$

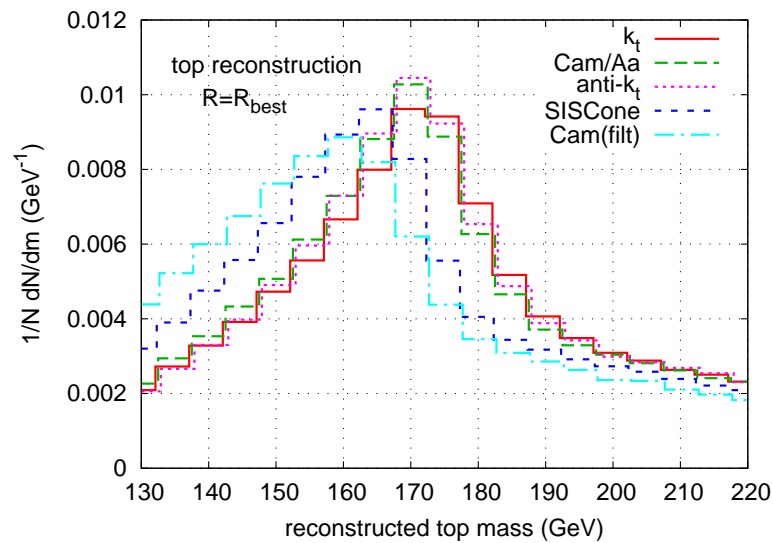
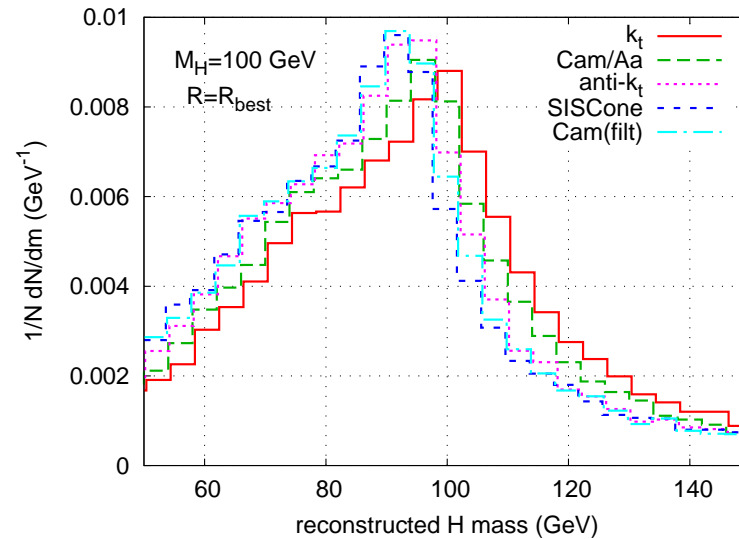
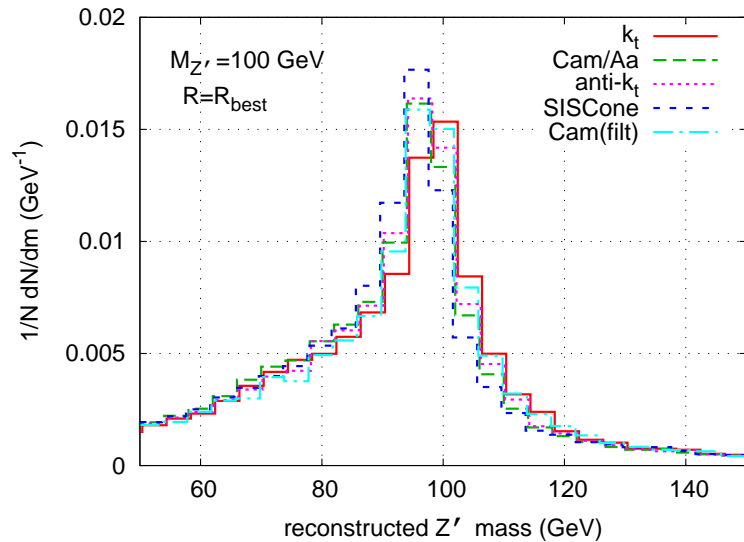
smaller width  $\equiv$  better signal-to-background ratio

- we can associate an **effective luminosity ratio**

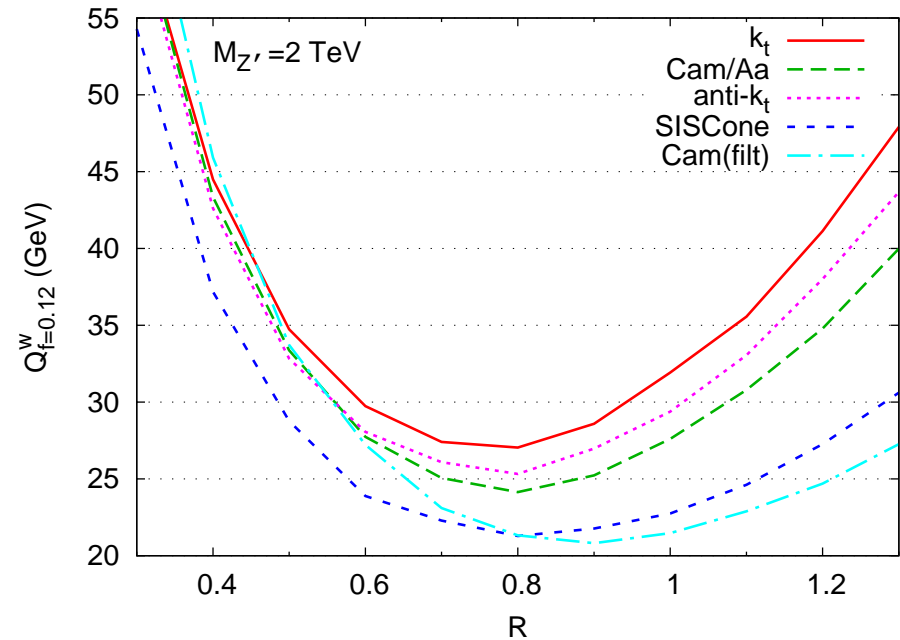
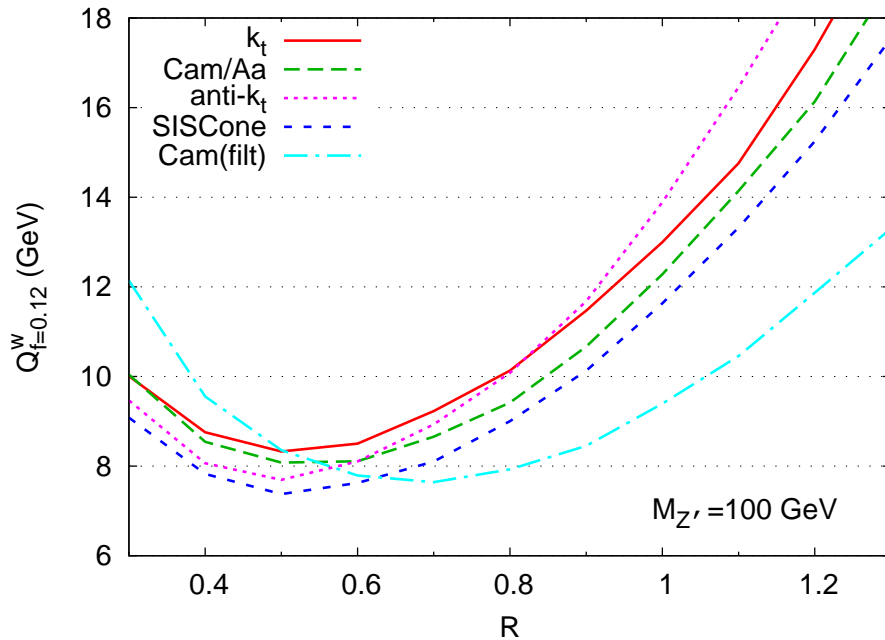
$$\rho_{\mathcal{L}} = \frac{\mathcal{L}_1}{\mathcal{L}_2} = \left[ \frac{(S/\sqrt{B})_{JA_1, R_1}}{(S/\sqrt{B})_{JA_2, R_2}} \right]^2 = \frac{Q_{f=z}^w(JA_2, R_2)}{Q_{f=z}^w(JA_1, R_1)}$$

e.g. if  $Q_{f=z}^w(JA_2, R_2) = 2Q_{f=z}^w(JA_1, R_1)$ ,  $(JA_2, R_2)$  will need twice the luminosity of  $(JA_1, R_1)$  to achieve the same discriminative power.

we see peaks...



**Message 1: there is a strong  $R$  dependence**



At 100 GeV,

using  $R = 0.8$  instead of  $R = 0.5$  means a discr. power loss of 20% ( $\rho_{\mathcal{L}} \approx 0.8$ )

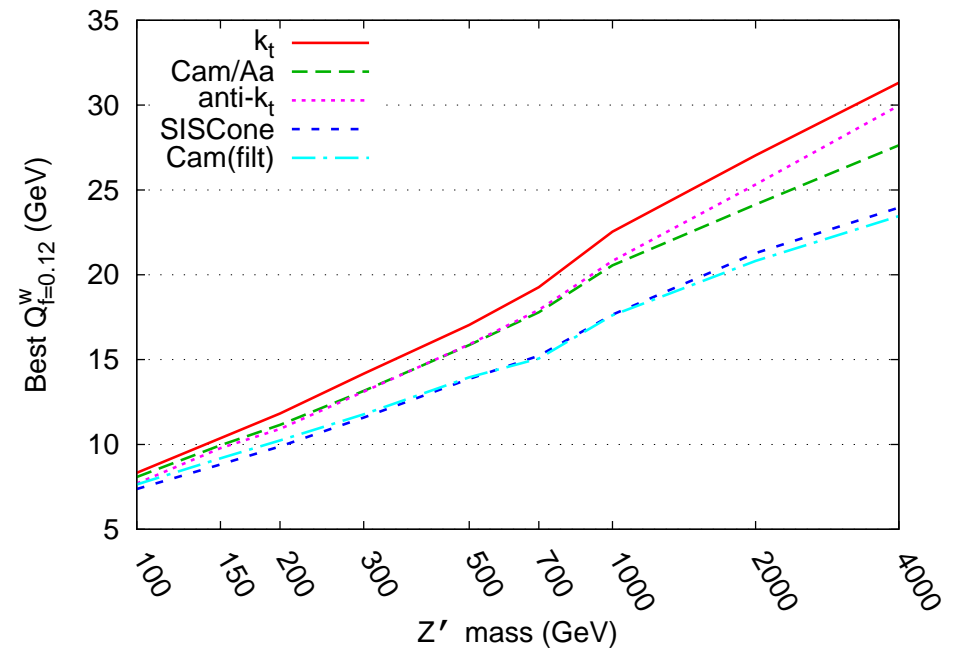
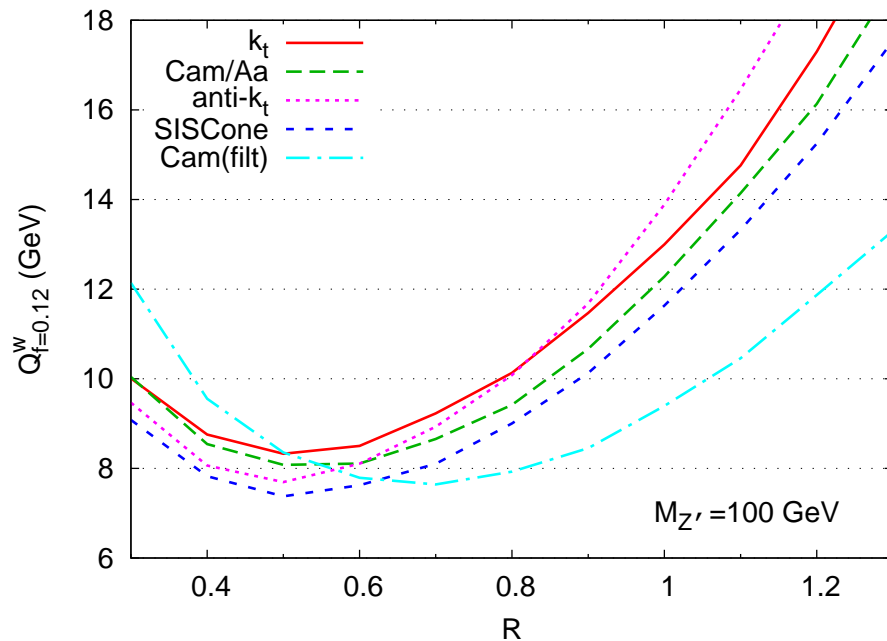
$R = 1$

40% ( $\rho_{\mathcal{L}} \approx 0.6$ )

At 2 TeV,

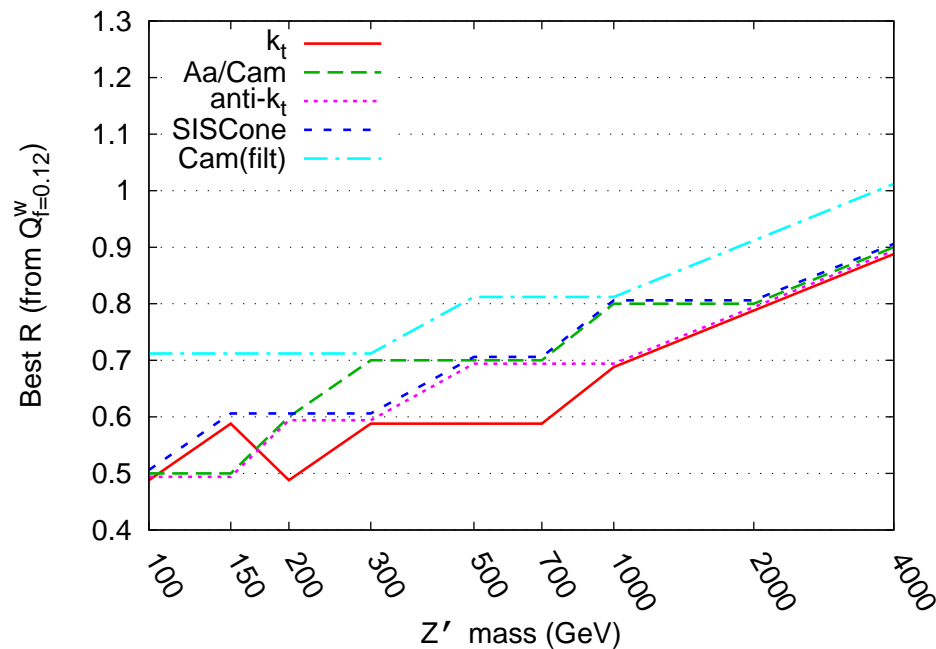
using  $R = 0.5$  instead of  $R = 0.8$  means a discr. power loss of 20% ( $\rho_{\mathcal{L}} \approx 0.8$ )

**Message 2: SIScone and Cam+filt do a slightly better job**



Using  $k_t$  instead of SIScone means a discr. power loss of  
 15% at 100 GeV ( $\rho_{\mathcal{L}} \approx 0.85$ )  
 20% at 2 TeV ( $\rho_{\mathcal{L}} \approx 0.8$ )

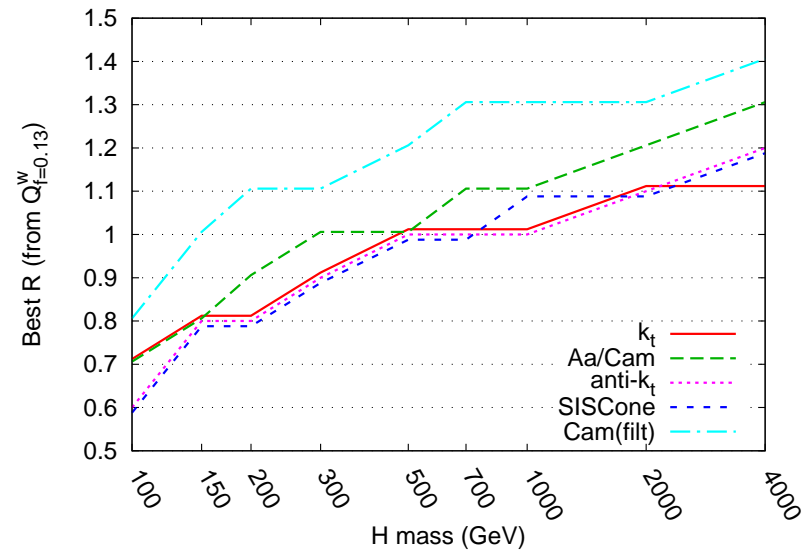
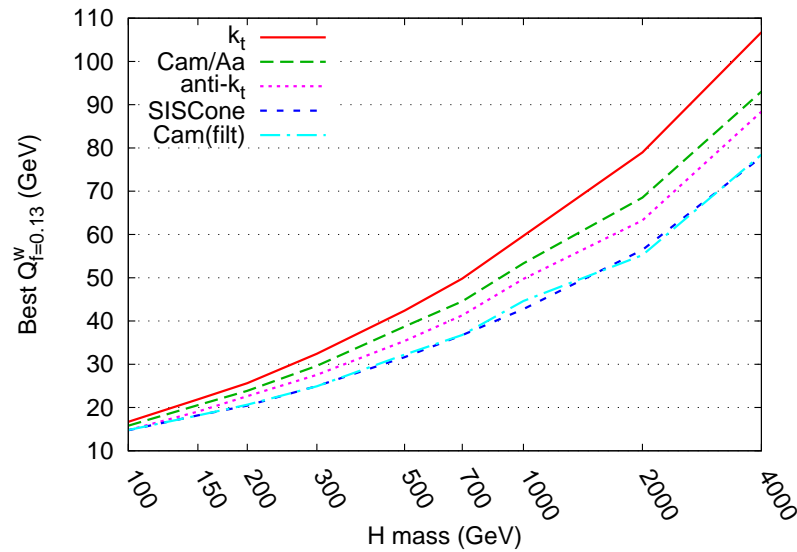
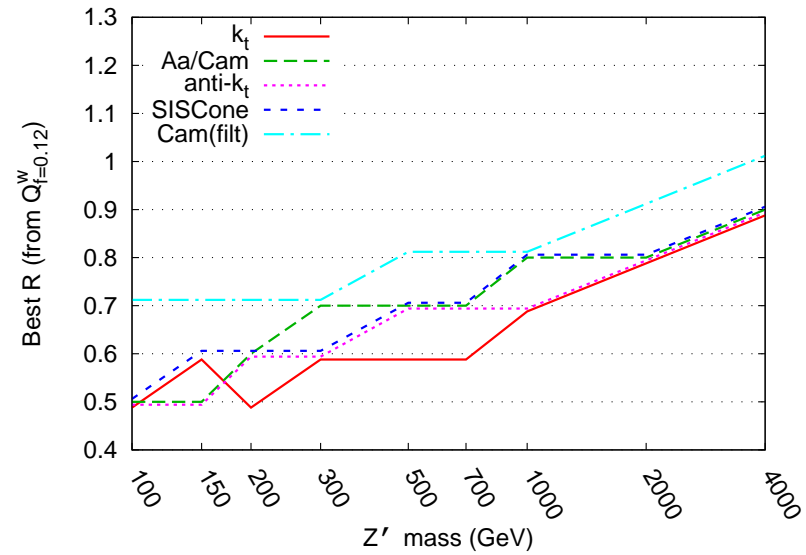
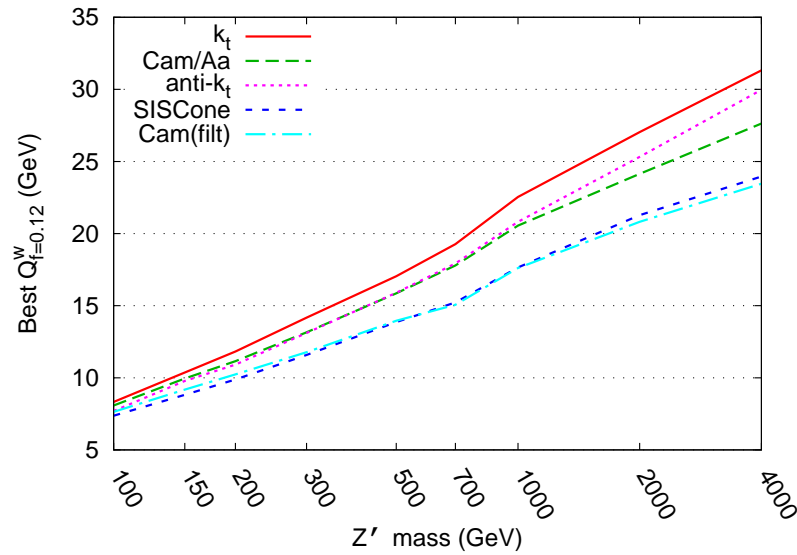
## Message 3: The parameters vary with the scale



The preferred value for  $R$  increase with the mass scale (typically like  $\log(M)$ )

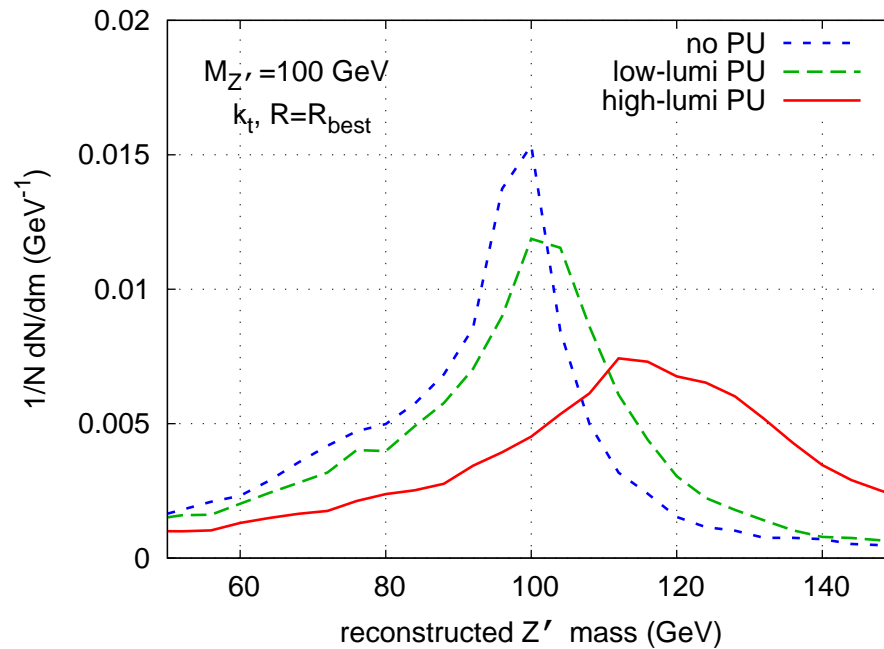


## Message 4: same for the gluon jets, though with a larger $R$



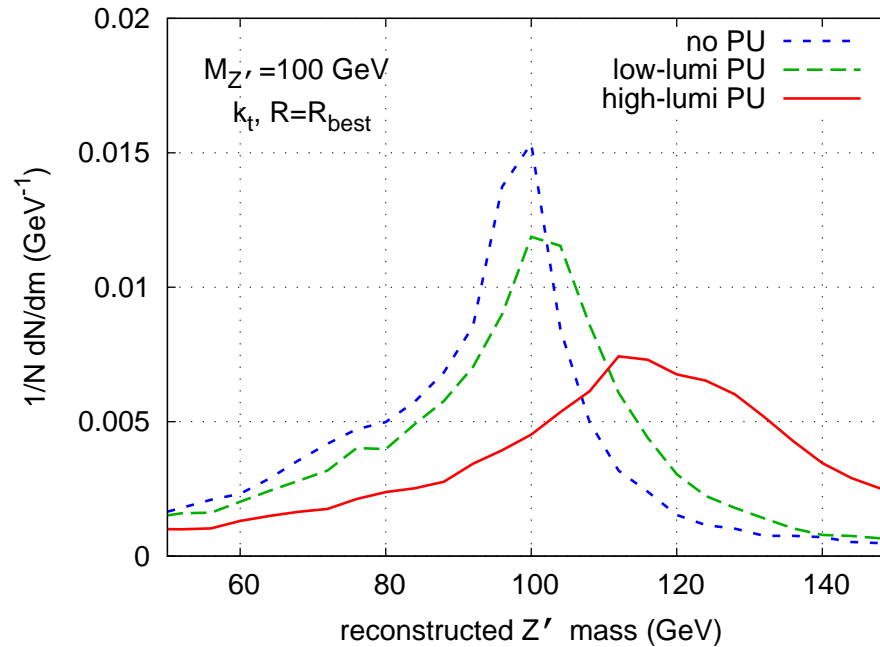
***Part 1***  
***Jets in  $pp$  collisions***  
***Subtracting pileup background using jet areas***

Pileup  $\approx$  uniform soft background that shifts jets to higher  $p_t$



... that needs to be subtracted!

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... that needs to be subtracted!

⇒ Using jet areas!

- Idea: add soft particles (**ghosts**) and look in which jets they are caught

jet area = region where it catches ghosts

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- 2 definitions
  - Passive area  
add one ghost and look where it ends. repeat to cover the  $(y, \phi)$  plane
  - Active area  
add a large amount of ghosts and cluster everything  
also gives purely ghosted jets

- Idea: add soft particles (**ghosts**) and look in which jets they are caught

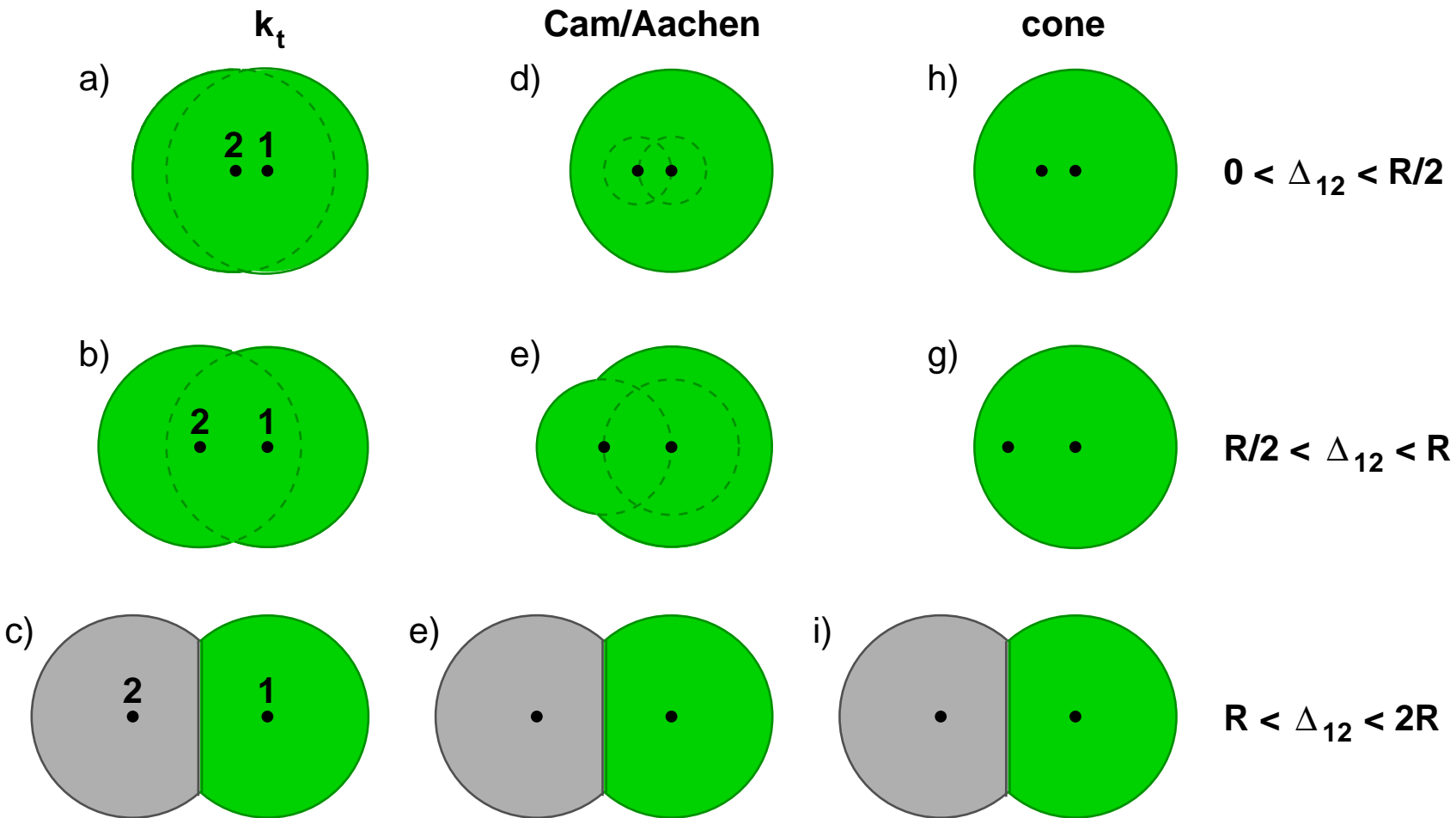
jet area = region where it catches ghosts

- 2 definitions
  - Passive area  
add one ghost and look where it ends. repeat to cover the  $(y, \phi)$  plane
  - Active area  
add a large amount of ghosts and cluster everything  
also gives purely ghosted jets
- Both definitions agree for dense events
- Both **practical** and **tractable analytically**

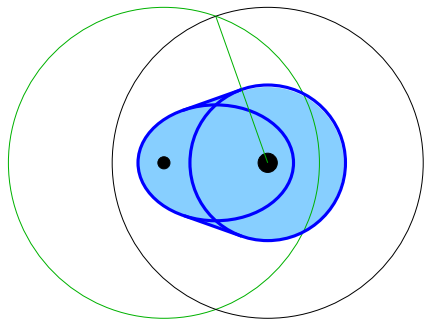
	$k_t$	Aac/Cam	cone
Passive	$\pi R^2$	$\pi R^2$	$\pi R^2$
Active	<p>(a) <math>k_t</math> algorithm</p> <p>pure ghost jets ———</p> <p>jets with 1 hard parton - - -</p> <p><math>\frac{A_{\text{hard}}}{\pi R^2} \approx 0.812 \pm 0.277</math></p> <p><math>\frac{A_{\text{ghost}}}{\pi R^2} \approx 0.554 \pm 0.174</math></p>	<p>(b) Cam/Aachen algorithm</p> <p>pure ghost jets ———</p> <p>jets with 1 hard parton - - -</p> <p><math>\frac{A_{\text{hard}}}{\pi R^2} \approx 0.814 \pm 0.261</math></p> <p><math>\frac{A_{\text{ghost}}}{\pi R^2} \approx 0.551 \pm 0.176</math></p>	<p><math>\frac{A_{\text{hard}}}{\pi R^2} = 0.25</math></p>



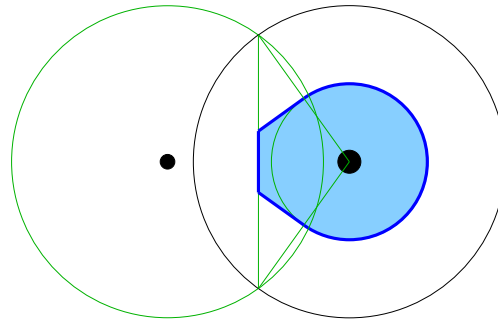
Passive area: 1 hard particle + 1 soft



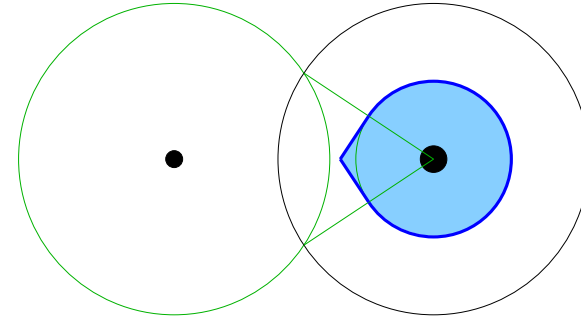
Active area: 1 hard particle + 1 soft: **analytic result for cone only**



$d < R$



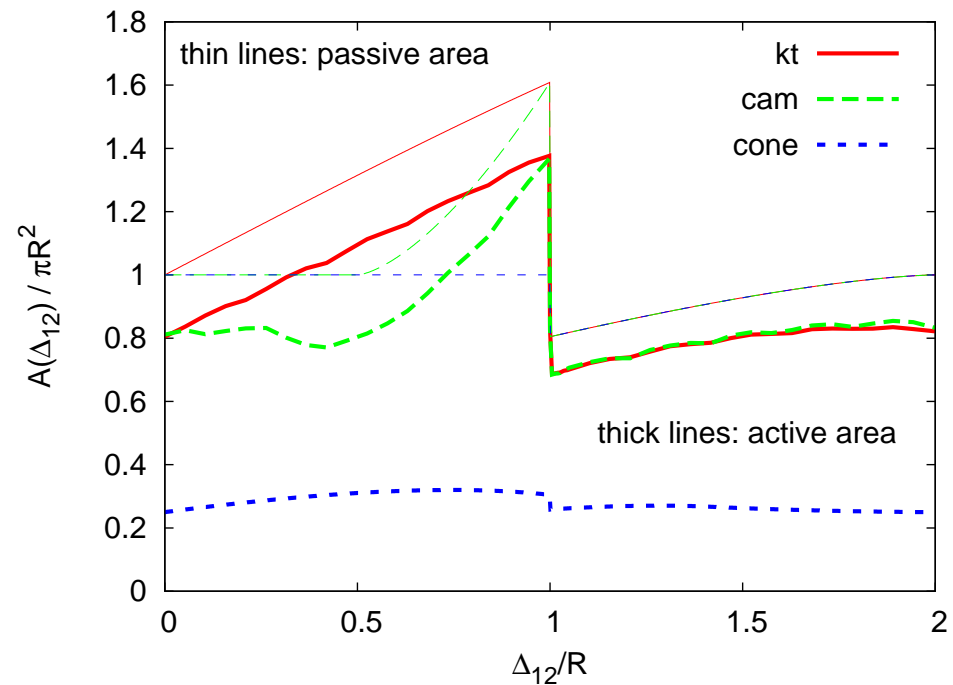
$R < d < \sqrt{2} R$



$\sqrt{2} R < d < 2R$

Alltogether, we have:

- Area  $\neq$  cst.  $\pi R^2$
- $\Delta_{12}$  dependence under control



QCD probability of emitting a small-angle soft gluon:

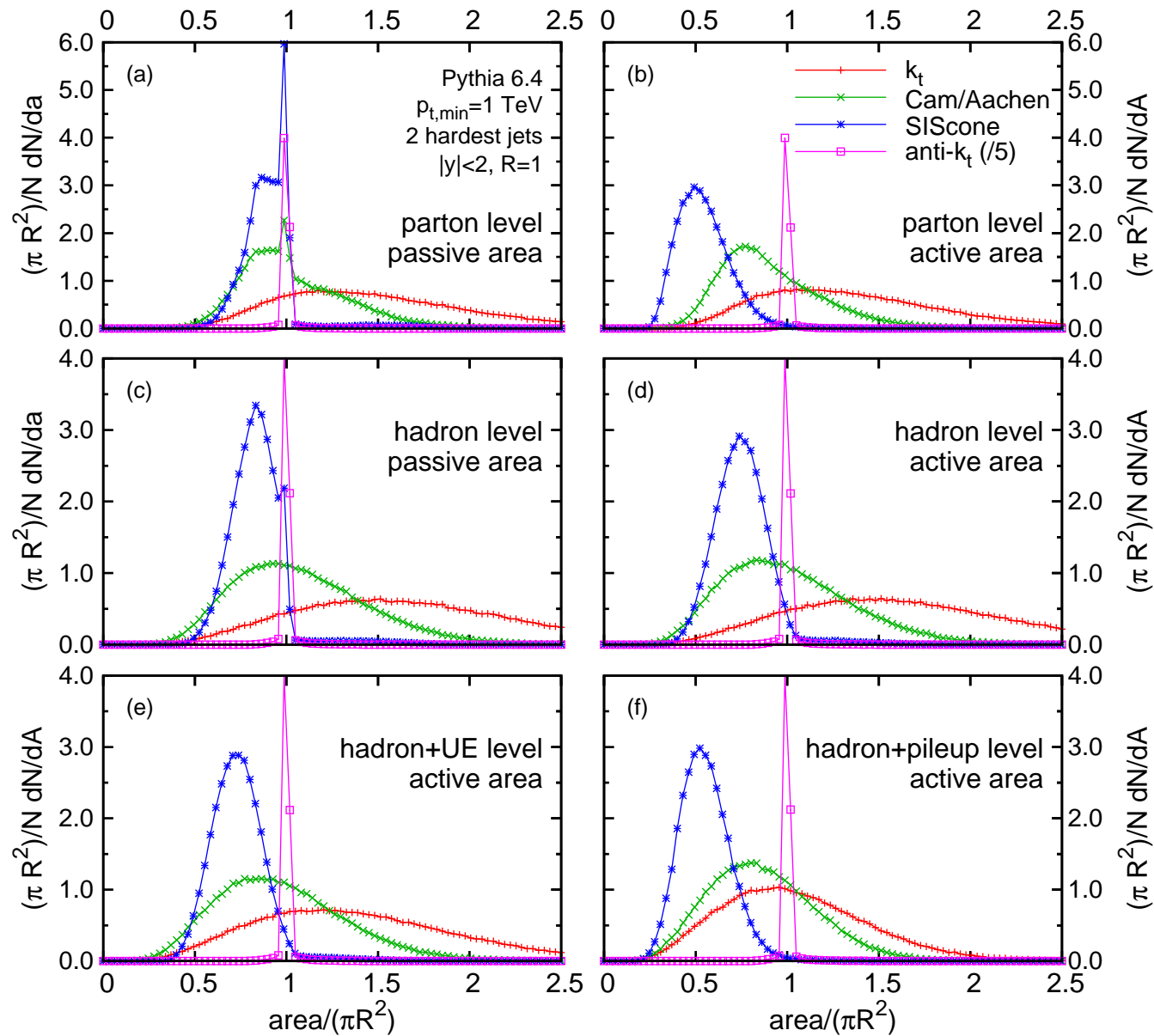
$$\frac{dP}{d\Delta_{12} dp_{t,2}} = C_{F,A} \frac{2\alpha_s}{\pi} \frac{1}{\Delta_{12}} \frac{1}{p_{t,2}}$$

Hence the average area is

$$\begin{aligned} \langle \mathcal{A}(p_{t,1}, R) \rangle &= \mathcal{A}_{\text{1hard}}(R) + \int d\Delta dp_{t,2} \frac{dP}{d\Delta_{12} dp_{t,2}} [\mathcal{A}_{\text{hard+1 soft}}(\Delta, R) - \pi R^2] \\ &= \frac{C_{F,A}}{\pi b_0} \log \left( \frac{\alpha_s(Q_0)}{\alpha_s(Rp_t)} \right) \pi R^2 d \end{aligned}$$

- Scaling violation
- gluon > quark
- with know LO anomalous dimension

$d$	passive	active
$k_t$	0.5638	0.519
Cam	0.07918	0.0865
SISCone	-0.06378	0.1246
anti- $k_t$	0	0



## Basic idea:

$$p_{t,\text{subtracted}} = p_{t,\text{jet}} - \rho_{\text{pileup}} \times \text{Area}_{\text{jet}}$$

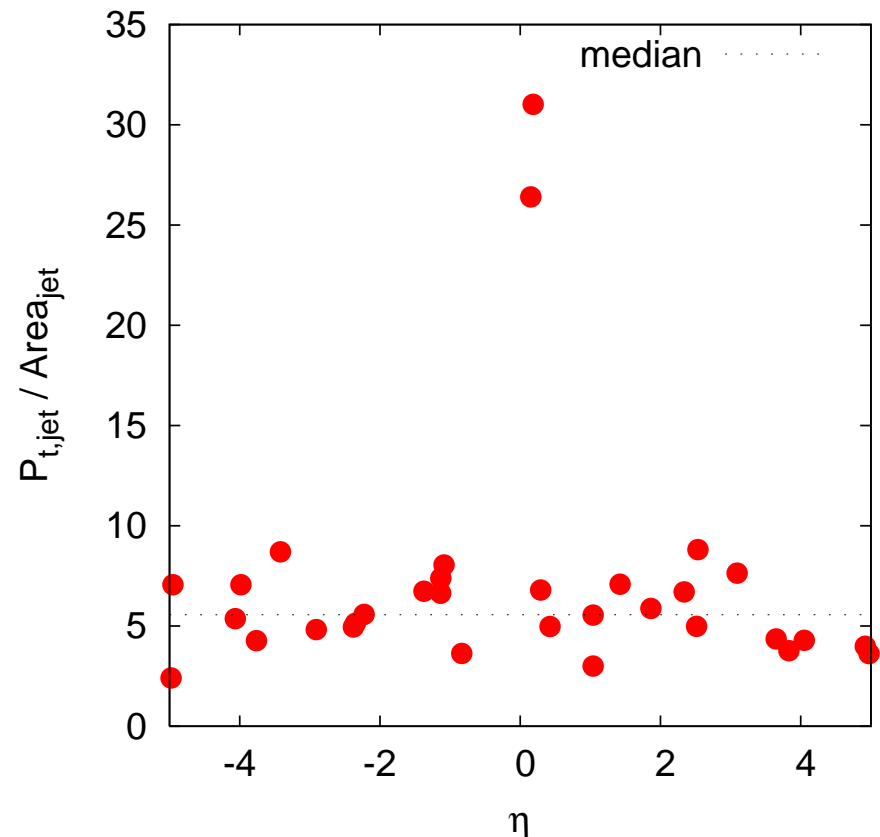
### ● Jet area:

- region where the jet catches infinitely soft particles
- tractable analytically

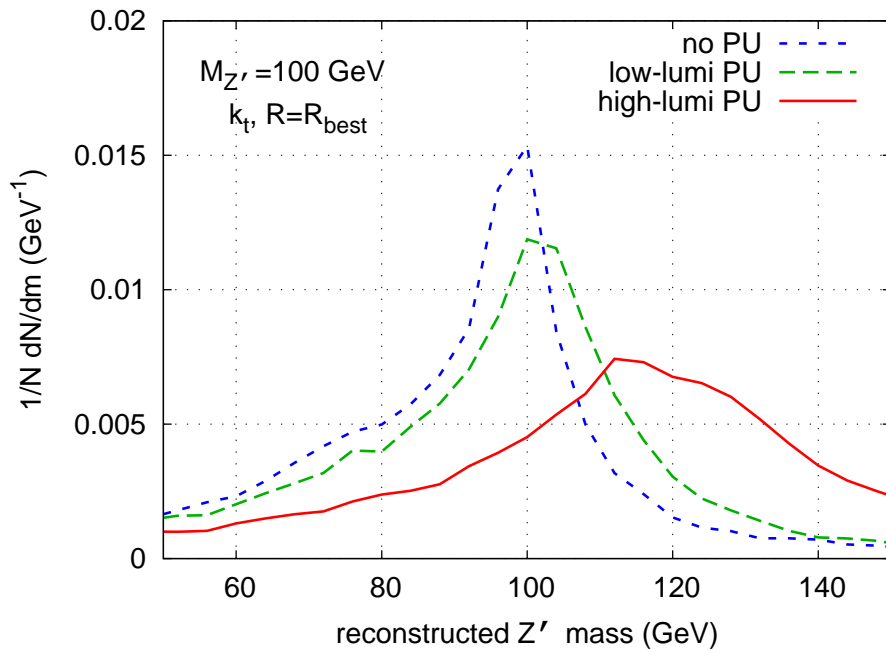
### ● Pileup density per unit area: $\rho_{\text{pileup}}$

e.g. estimated from the median  
of  $p_{t,\text{jet}} / \text{Area}_{\text{jet}}$

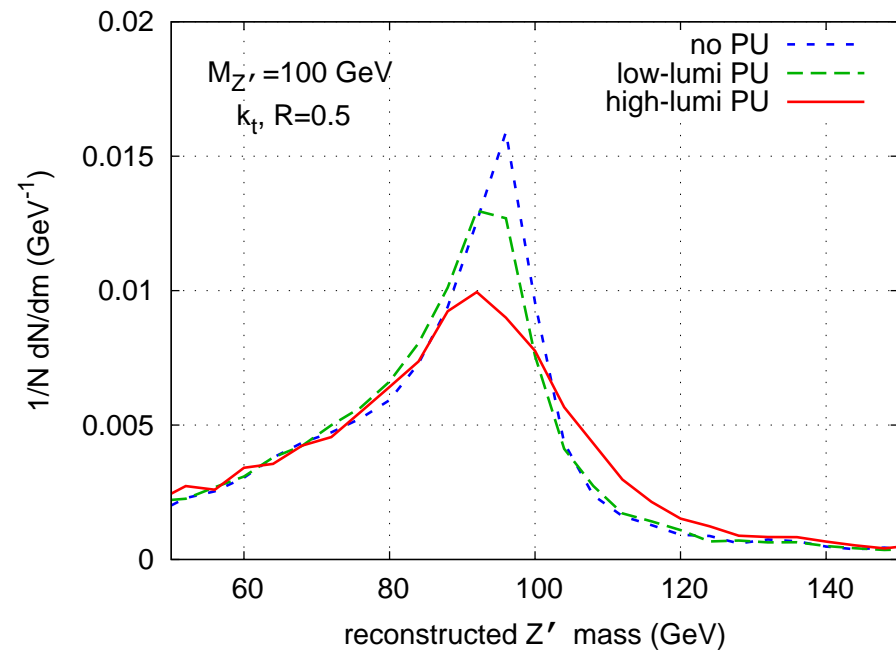
implemented in FastJet  
on an event-by-event basis



# Subtraction recipe

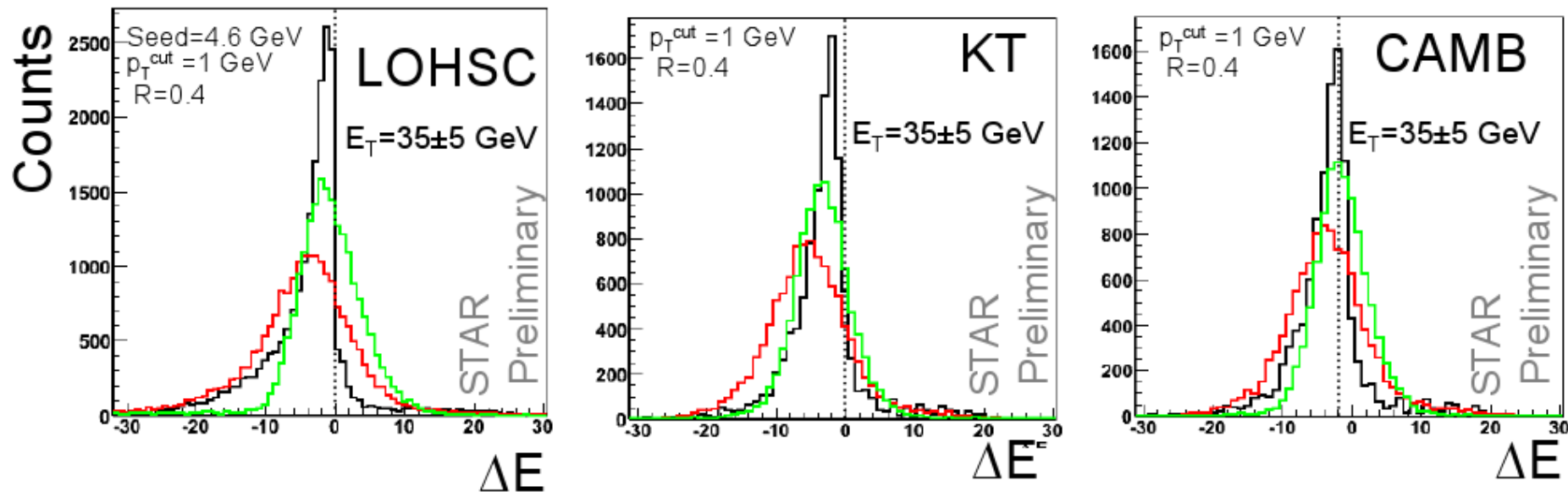


subtraction



***Part 3: heavy-ion background:  
subtracting more complex background***

[S. Salur, J. Putschke, ... (STAR)]



Red curves:  $\Delta E = E^{\text{PyEmbed}} - E^{\text{PyDet}}$

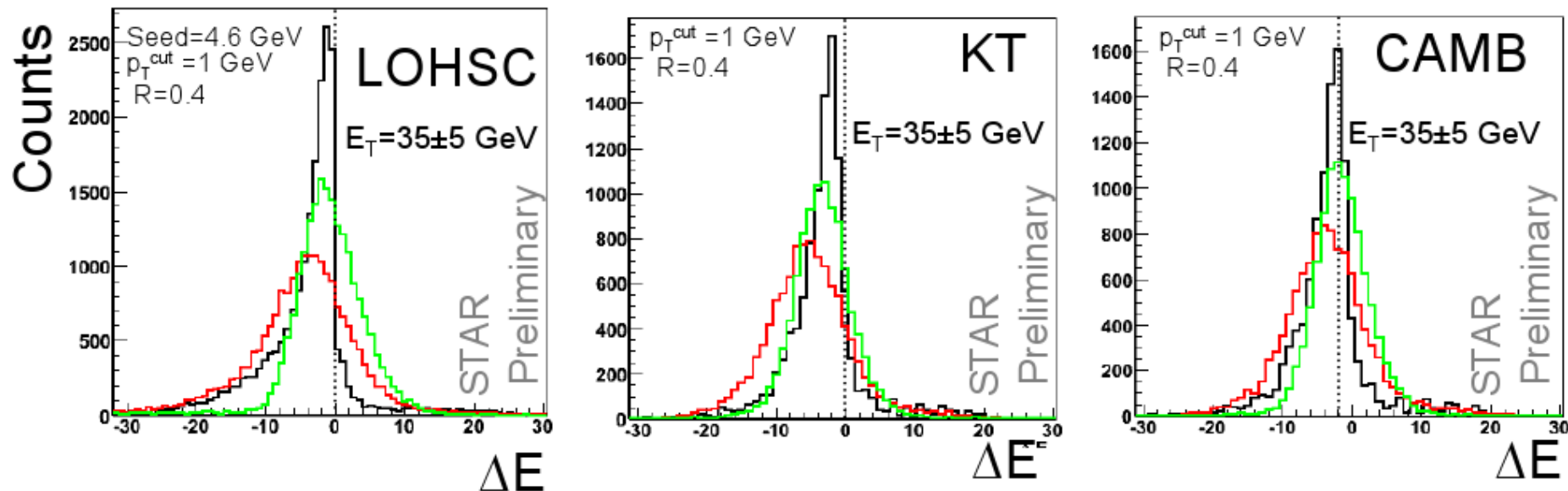
with  $\text{PyDet} \equiv \text{Pythia } pp + \text{detector effects}$

$\text{PyEmbed} \equiv \text{same with real AuAu event added}$

⇒ measure of the subtraction efficiency



[S. Salur, J. Putschke, ... (STAR)]



Red curves:  $\Delta E = E^{\text{PyEmbed}} - E^{\text{PyDet}}$

with PyDet  $\equiv$  Pythia  $pp$  + detector effects

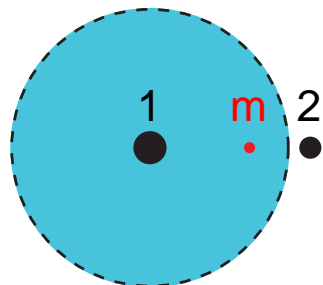
PyEmbed  $\equiv$  same with real AuAu event added

$\Rightarrow$  measure of the subtraction efficiency

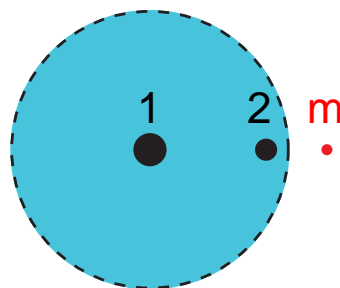
Work under progress: removing the last few GeV shift in  $\Delta E$

Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
- Modify the hard scattering (back-reaction)
  - can be pointlike or diffuse
  - gain:  $p_2$  gained when adding  $p_m$

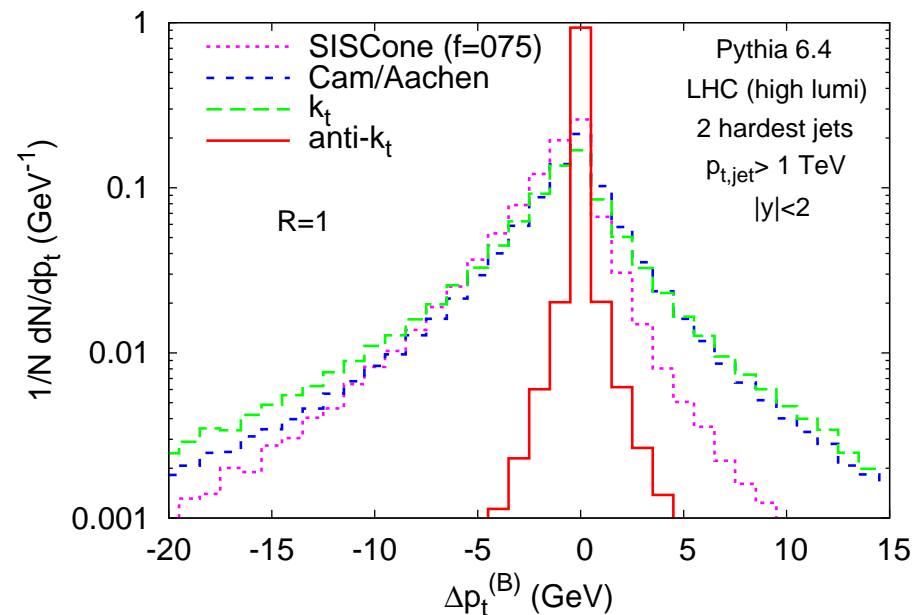


- loss:  $p_2$  lost when adding  $p_m$



Additional soft background has 2 effects:

- **Throw soft particles in the hard jet:** dealt with by subtraction
- **Modify the hard scattering (back-reaction)**
  - can be pointlike or diffuse
  - tractable analytically (similar to areas)
  - $k_t \gtrsim$  Cambridge  $>$  SIScone  $\gg$  anti- $k_t$

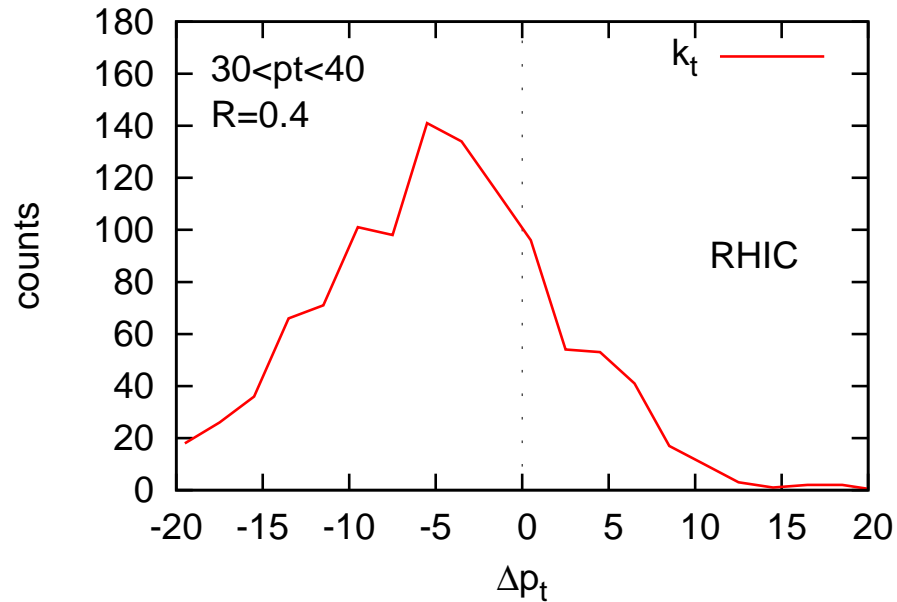


Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
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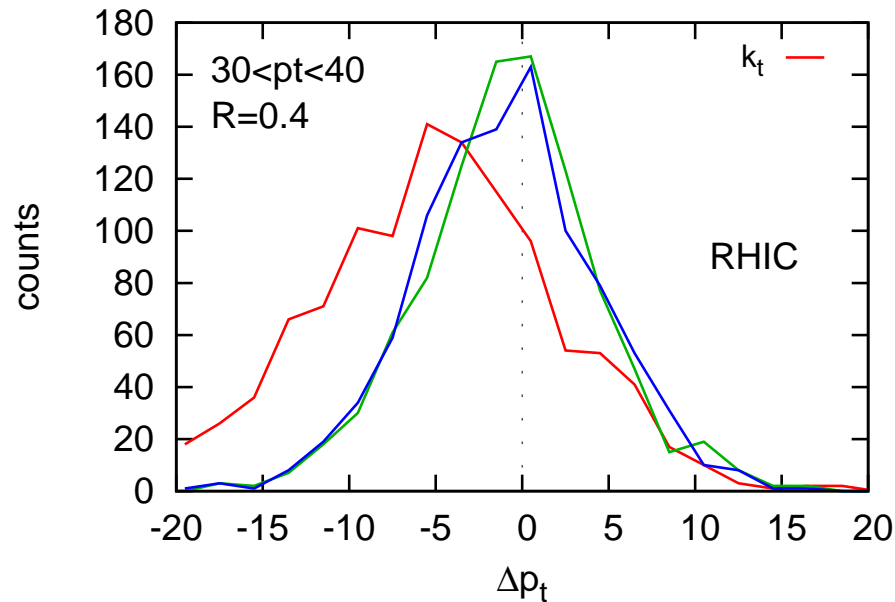
+ For heavy-ion collisions: fluctuating underlying event background  
→ median estimation of  $\rho$  might oversimplified

Work under progress: test subtraction using  
Pythia  $pp$  hard event + HYDJET  $AA$  background



Similar shifts than STAR

Work under progress: test subtraction using  
Pythia  $pp$  hard event + HYDJET  $AA$  background



Target: reach a precision of 1-2 GeV by

- carefully tuning the algorithm to reduce UE sensitivity and back-reaction
- carefully tuning the subtraction to deal with the fluctuating background

- Use IRC-safe algorithms!
- Jet-finding in  $pp$  at the LHC
  - SISCone and Cam+filt. do a slightly better job
  - strong  $R$  dependence: important to choose  $R_{\text{best}}$
  - $R_{\text{best}}$  increases with the scale
  - same for quark and gluon jets, larger  $R_{\text{best}}$  for gluons

⇒ flexibility in jet finding at the LHC
- Subtraction using jet areas
  - Jet areas: clearly defined, analytic control
  - Simple systematic pileup subtraction
- Jet-finding in  $AA$  at RHIC and the LHC
  - First measurement at RHIC
  - Work under progress: improve subtraction down to 1-2 GeV