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Twistor techniques in QCD

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Motivation

Calculation of QCD amplitudes

- Color decomposition
- Spinor-helicity formalism
- Twistor techniques
 - Tree-level techniques
 - One-loop techniques
- Conclusion

Motivation

- Most of the events observed at the LHC will be QCD events.
- Many interesting signatures correspond to a multijet final state, e.g.
 - $t \,\overline{t} \, H \rightarrow \ell \nu + 6$ jets, **8 jets** - $p \, p \rightarrow \tilde{q} \, \tilde{q} \rightarrow (q \, \tilde{\chi}^0) (q \, \tilde{\chi}^0) \rightarrow (q \, cds) (q \, cds)$ $p \, p \rightarrow 8$ jets



Motivation

- The background to the interesting LHC signatures is in very often given by high-multiplicity QCD events.
- Precision measurements require theory predictions beyond leading order.

$$\sigma = \underbrace{\alpha_s^n \sigma_0}_{\text{LO}} + \underbrace{\alpha_s^{n+1} \sigma_1}_{\text{NLO}} + \mathcal{O}(\alpha_s^{n+2})$$





Feynman diagrams

 Standard technique for the calculation of amplitudes in quantum field theory

Feynman diagrams

• Example:
$$e^+ e^- \rightarrow \gamma \gamma \gamma$$

6 diagrams at tree level.





Feynman diagrams

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• Example of a QCD process $gg \rightarrow ggg$ 45 Feynman diagrams at tree level.







Feynman diagrams

- The number of Feynman diagrams is growing extremely fast with the number of external particles.
 - 6 external gluons: 510
 - 7 external gluons: 5040
 - 8 external gluons: 40320
 - ...
- The reason for this fast growth in complexity is color...

Color decomposition

• Solution: Factor out the color information!

 $\mathcal{A}_{n}^{\text{tree}}\left(\{k_{i},\lambda_{i},a_{i}\}\right) = g^{n-2} \sum_{\sigma \in S_{n}/Z_{n}} \text{Tr}\left(T^{a_{\sigma(1)}} \cdots T^{a_{\sigma(n)}}\right) A_{n}^{\text{tree}}(\sigma(1^{\lambda_{1}}),\ldots,\sigma(n^{\lambda_{n}}))$

- The color-ordered amplitudes
 - correspond to a specific color-ordering.
 - depend only on the momenta and helicities of the gluons.
 (but not on color!)
- This construction can be easily extended to quarks as well as beyond tree-level.

Spinor-helicity formalism

- As the color-ordered amplitudes depend only on the momenta and helicities, they can be easily calculated using spinor-helicity formalism.
- Example: The color-ordered amplitude $g g \rightarrow g g g$ for a specific helicity assignment:

$$A_5^{tree}(1^-, 2^-, 3^+, 4^+, 5^+) = \frac{\langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle}$$

where $\langle ij \rangle \sim \sqrt{(p_i + p_j)^2}$.

Spinor-helicity formalism

- The color-ordered amplitudes can be calculated via:
 - Color-ordered Feynman rules.
 - Recursive techniques (Berends-Giele recursion)
 - The recently introduced twistor techniques...





Twistors in QCD

- In 1967, R. Penrose introduces a new mathematical object called a twistor.
- In 2003, E. Witten conjectures a duality between QCD and a certain type of string theory based on Pensore's twistor space.
- Main message:

"Be complex, but stay on-shell!"





Twistors in QCD

- Tree-level:
 - BCFW recursion
 - CSW formalism
- One-loop:
 - Generalized unitarity





- First conjectured by Britto, Cachazo and Feng, proven analytically by Witten. [hep-th/0412308, hep-th/0501052]
- Main idea: QCD amplitudes can be build recursively from smaller, on-shell amplitudes.

$$A_n(1^+, 2, \dots, n^-) = \sum_{k=2}^{n-2} A_{k+1} \left(\hat{1}, 2, \dots, k, -\hat{P}_{1,k}^{-h} \right) \frac{1}{P_{1,k}^2} A_{n-k+1} \left(\hat{P}_{1,k}^h, k+1, \dots, \hat{n} \right)$$

















• Can be generalized to quarks.

[Schwinn, Weinzierl, hep-ph/0703021]

- Does not only hold at the level of the color-ordered amplitudes, but can be generalized to the full color-dressed QCD amplitudes.
 [CD, Hoesche, Maltoni, hep-ph/0607057]
- Holds also for QED amplitudes

[Ozeren, Stirling, hep-th/0509063]

 Very compact results for specific helicity configurations valid for an arbitrary number of gluons.





BCFW recursion

• Example of a 7 gluon amplitude from the BCFW recursion

$$\begin{aligned} A(1^{-}, 2^{-}, 3^{-}, 4^{+}, 5^{+}, 6^{+}, 7^{+}) &= \\ \frac{\langle 1|2+3|4|^{3}}{t_{2}^{[3]}\langle 5|6\rangle\langle 6|7\rangle\langle 7|1\rangle[2|3][3|4]\langle 5|4+3|2]} \\ &- \frac{1}{\langle 3|4\rangle\langle 4|5\rangle\langle 6|7+1|2]} \left(\frac{\langle 3|(4+5)(6+7)|1\rangle^{3}}{t_{3}^{[3]}t_{6}^{[3]}\langle 6|7\rangle\langle 7|1\rangle\langle 5|4+3|2]} + \frac{\langle 3|2+1|7|^{3}}{t_{7}^{[3]}\langle 6|5\rangle[7|1][1|2]} \right) \end{aligned}$$





 Any tree-level amplitude can be written as a sum of CSW diagrams
 [Cachazo, Svrcek, Witten, hep-th/0403042]

 $A_{6}(1^{-}, 2^{-}, 3^{-}, 4^{+}, 5^{+}, 6^{+}) =$

• 3 CSW diagrams vs. 510 Feynman diagrams.



CSW formalism

• Any tree-level amplitude can be written as a sum of CSW diagrams [Cachazo, Svrcek, Witten, hep-th/0403042]

 $A_6(1^-, 2^-, 3^-, 4^+, 5^+, 6^+) =$



The vertices are on-shell QCD amplitudes (for a specific helicity assignment) of the form

$$\frac{\langle 1P \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 4P \rangle \langle P1 \rangle}$$



CSW formalism

• Any tree-level amplitude can be written as a sum of CSW diagrams [Cachazo, Svrcek, Witten, hep-th/0403042]

 $A_6(1^-, 2^-, 3^-, 4^+, 5^+, 6^+) =$



 $\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle$

 $(P)(P1) \qquad Complex on-shell momentum$



Can be generalized to quarks, and scalars.

[Georgiou, Khoze, hep-th/0404072; Badger, Glover, Khoze, hep-th/0412275]

- Very easy diagrammatic technique, giving very compact results.
- Suitable not only for the calculation of amplitudes, but also for the calculation of
 - splitting functions
 - antenna functions

[Birthwright, Glover, Khoze, Marquard, hep-ph/0503063]

[CD, Maltoni]



Generalized unitarity

 All one-loop QCD amplitudes can be reduced to a sum of boxes, triangles and bubbles:



 If we cut one line (i.e. put on-shell one loop propagator), unitarity tells us how the amplitude behaves:



[Bern, Dixon, Dunbar, Kosower, hep-ph/0943226]

 $\langle \neg \rangle$



• All one-loop QCD amplitudes can be reduced to a sum of

boxes, triangles and bubbles:



• If we cut a second line... the result is zero!

But this assumes = 0! the momenta to be real...





Generalized unitarity

In complex momenta, more than one unitarity cut is allowed:



Generalized unitarity

In complex momenta, more than one unitarity cut is [Britto, Cachazo, Feng, hep-th/0412103; allowed:

Anastasiou, Britto, Feng, Kunzt, Mastrolia, hep-ph/0607011; Ossola, Papadopulos, Pittau, hep-ph/0802.1876]



 $\mathbf{c}_{4} \quad = \mathcal{A}^{tree} \mathcal{A$

- The tree amplitudes can be calculate using other techniques (CSW, BCFW, ...).
- This fixes completely the coefficient C4.



Conclusion

- Feynman diagrams are not the most efficient technique to calculate gauge theory amplitudes.
- Color-ordered amplitudes can be very easily and efficiently calculated using the twistor techniques, both at tree-level and at one-loop level.
- The results obtained from these techniques have in general very simple and compact analytic expressions.













• Par exemple pour $gg \rightarrow ggg$, on a

Digarammes de Feynman	Diagrammes CSW





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45	





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